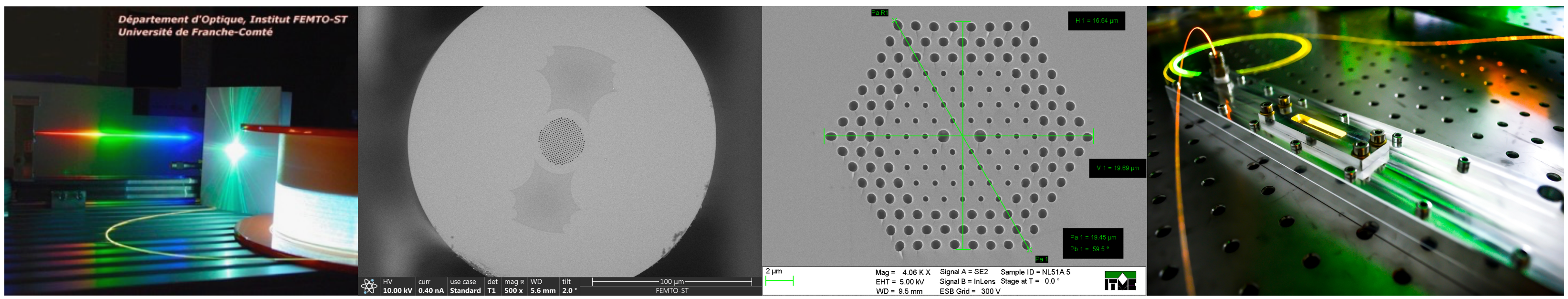
femto-st SCIENCES & TECHNOLOGIES

Nonlinear Effects in Optical Fibers : From Fundamentals to Applications



Smart Systems for a better life - SMYLE

Improving French-Swiss partnerships in engineering

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UBFC UNIVERSITÉ BOURGOGNE FRANCHE-COMTÉ

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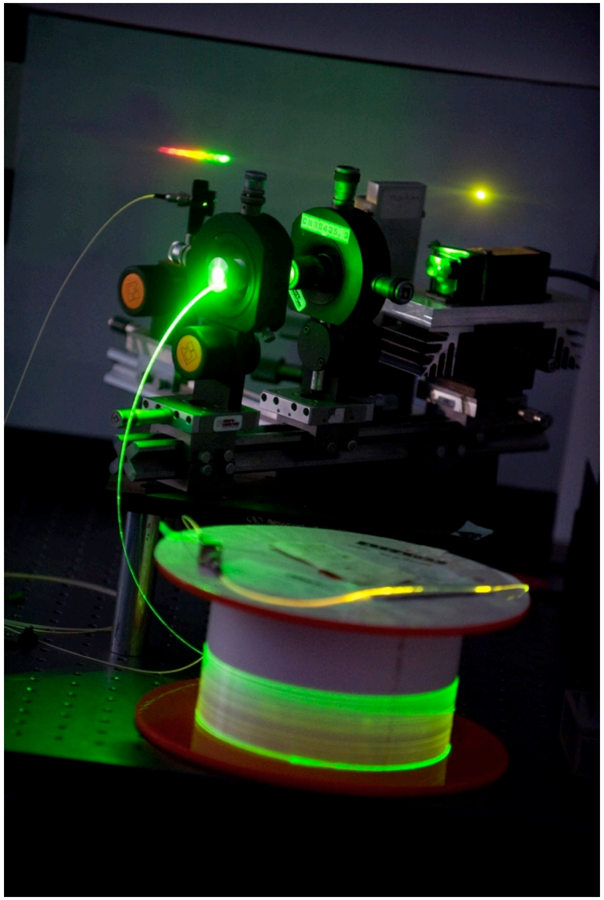
Outline

Contents



Introduction to Nonlinear Optics (NLO)

1672 : **The Newton's Experiment**

Three centuries ago, Newton discovered that white light is the superposition of coloured rays.

Then he succeeded in reconstructing white light using a second prism.

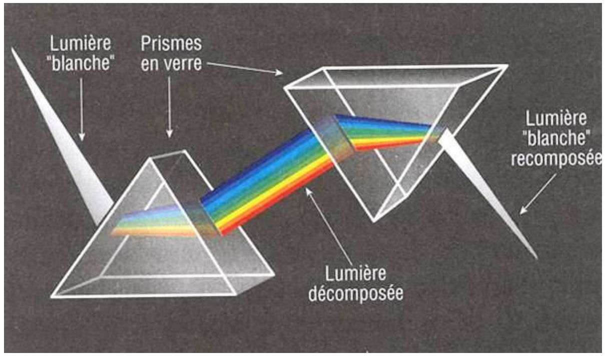
He then tried to alter the color of pure rays, and he tells….

"They obstinately retained their colours, notwithstanding the utmost endeavours to change it…."

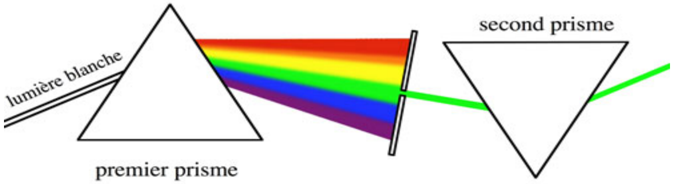
Isaac Newton (1643-1727)

« Newton's Experimentum Crucis »

« White light reconstruction using a second prism »



«Newton' quest : change the color of a pure ray »



Introduction to Nonlinear Optics (NLO)

- 1960 : **Discovery of the LASER,**
  high- brightness, coherence, monochromaticity, enabling strong light-matter interactions

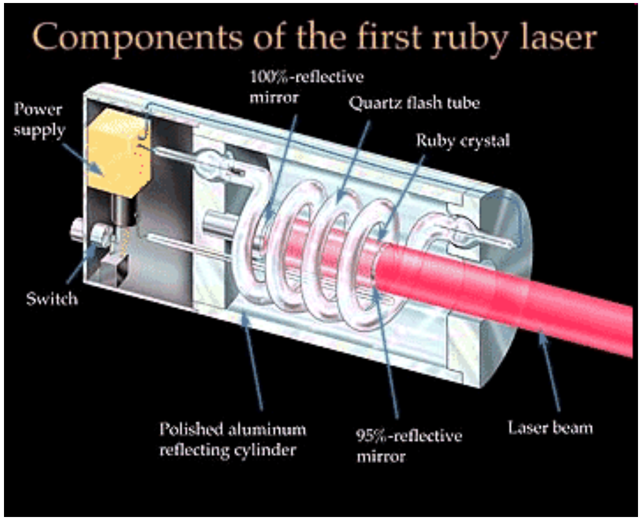
- 1961 : **Modern Nonlinear Optics (NLO)**,
  Second harmonic generation (SHG)

- 1965-1990: Stimulated Raman scattering (SRS) Third harmonic generation (THG), Difference and Sum frequency generation (DFG, SFG), Parametric amplification (OPA), Parametric oscillation (OPO), Four-wave mixing (FWM), Phase conjugation (PC), self-Focusing, Modulation instability, Solitons, Supercontinuum generation

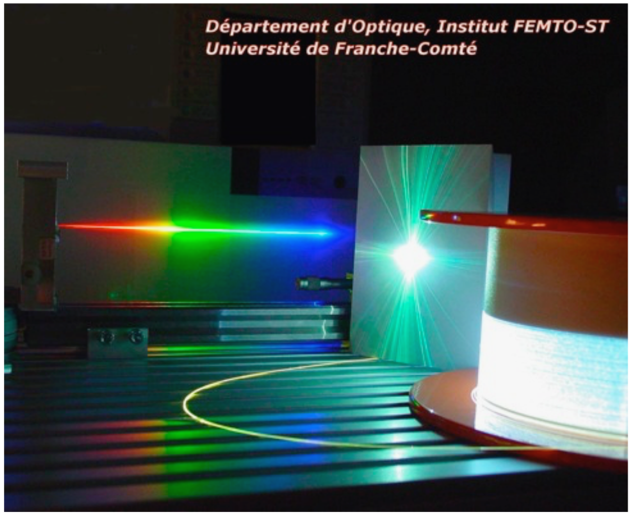
- 2000 : **Supercontinuum generation in photonic crystal fiber = The « Rainbow » laser:**
  Generation of a spatially-coherent white light from a photonic crystal fiber

Definition: « **In nonlinear optics, we are concerned with the effects that the light induces on itself as it propagates through the medium** »

« The first Ruby laser » by Maiman

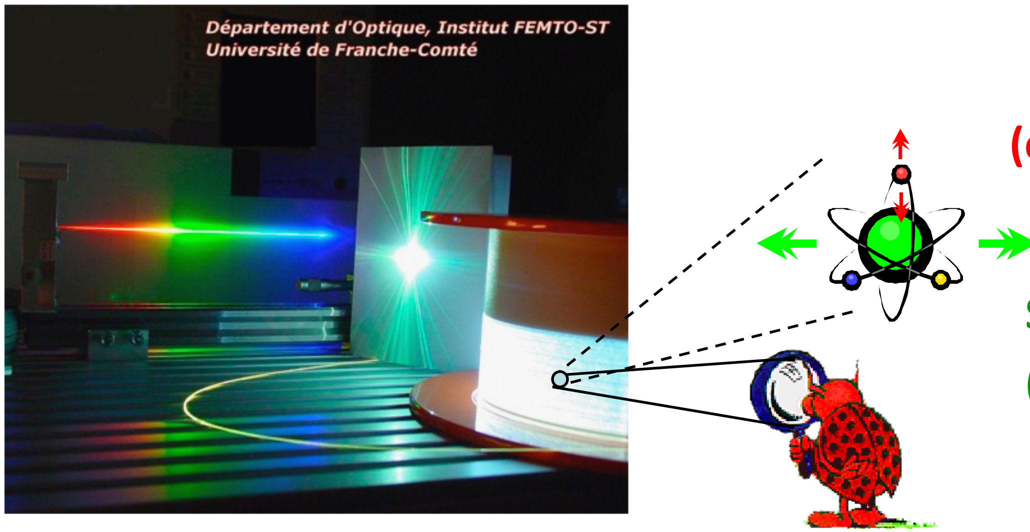


« The white-light laser » in 2000



Third-order Nonlinear Effects

Nonlinear optical effects are remarkable phenomena that arise when an intense optical beam propagates through an optical fiber



Optical Kerr Effect
(electronic response)
$n=n_0+n_2I(z,t)$

Stimulated Raman Scattering
(Delayed Molecular response)

Third-Order Nonlinear Polarization

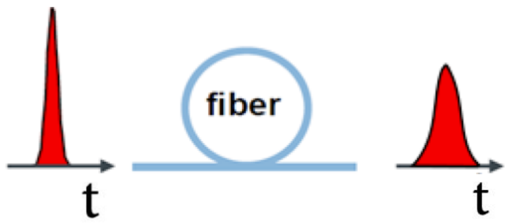
$$P_{NL}(t)=\varepsilon_0\chi_K^{(3)}:E(t)E(t)E(t)+\varepsilon_0 E(t)\int_{-\infty}^{t}\chi_R^{(3)}(t-t')E(t')E(t')dt'$$

Optical Kerr effect:
Instantaneous Elastic effect
(No energy exchange with matter)

Raman effect :
Inelastic light scattering:
(Molecular vibrational states)

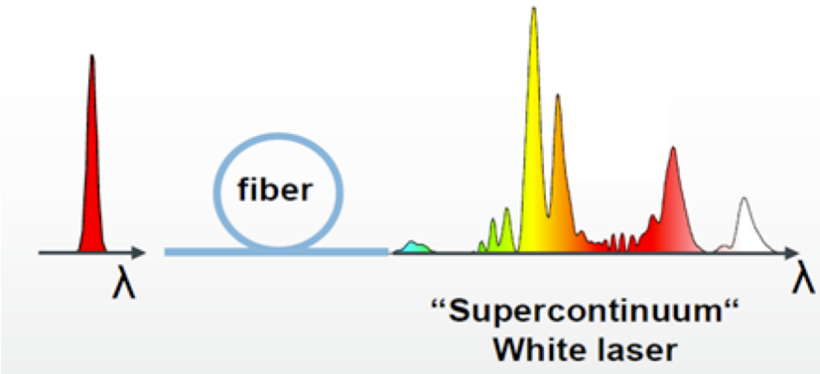
Linear and nonlinear propagation in optical fibers

- **Linear effects : Dispersion & Absorption**



- **Nonlinear effects :**

  Self-phase modulation (SPM)
  Optical Solitons (OS),
  Stimulated Raman Scattering (SRS)
  Dispersive Waves (DW)
  Optical Wave Breaking (OWB)
  Four-Wave Mixing (FWM)
  Modulation Instability (MI)
  All intensity dependent



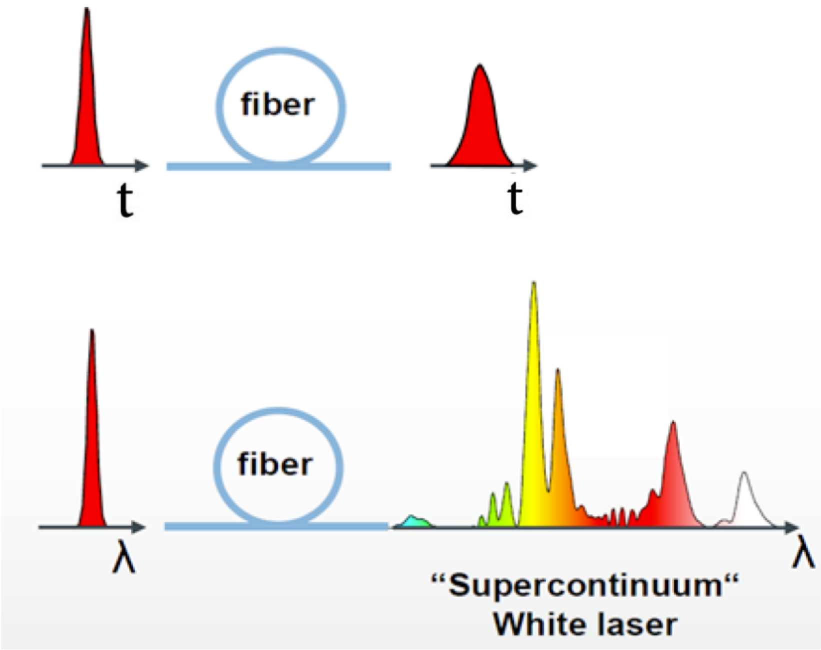
- SC is broad as the sun (from UV to IR) and bright as a laser > 20k the sun
- Spatially coherent - single-mode beam output - Fiber delivery

- **Supercontinuum light sources have many applications: OCT, Absorption Spectroscopy, Microscopy, Biomedical Imaging, OFC Metrology, etc...**

R. R. Alfano, The Supercontinuum Laser Source (Springer, New York, 2016).

Linear and nonlinear propagation in optical fibers

- **Linear effects : Dispersion & Absorption**



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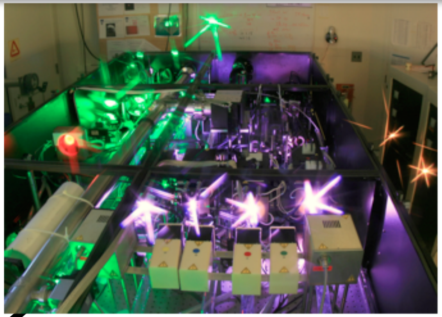
Key Applications of Nonlinear Optics

Telecommunications

Fiber Sensors



Fiber Lasers

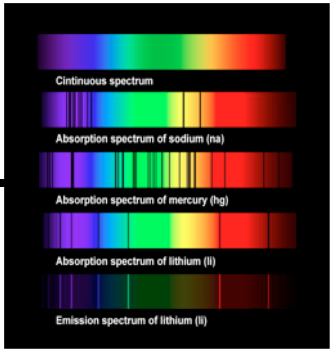
Extreme & Quantum physics

Kerr

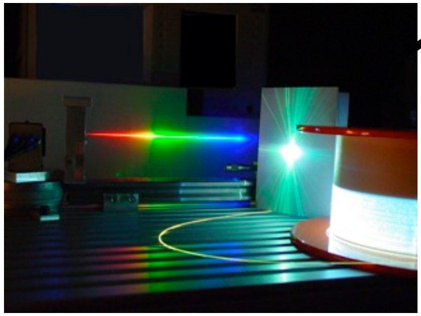
Raman

Brillouin

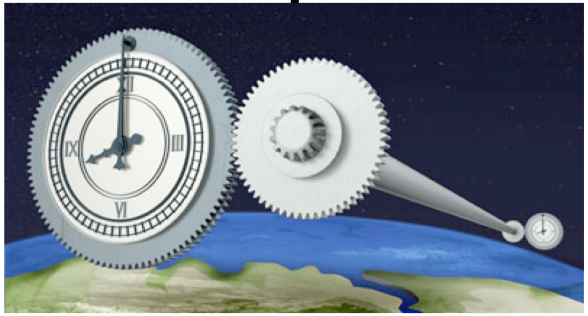
Nonlinear effects



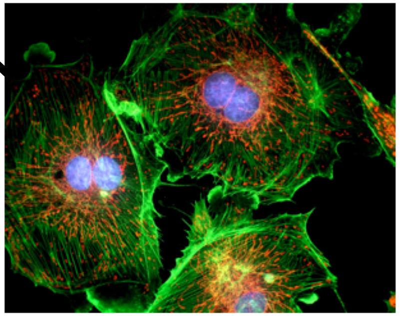
Spectroscopy



Supercontinuum sources



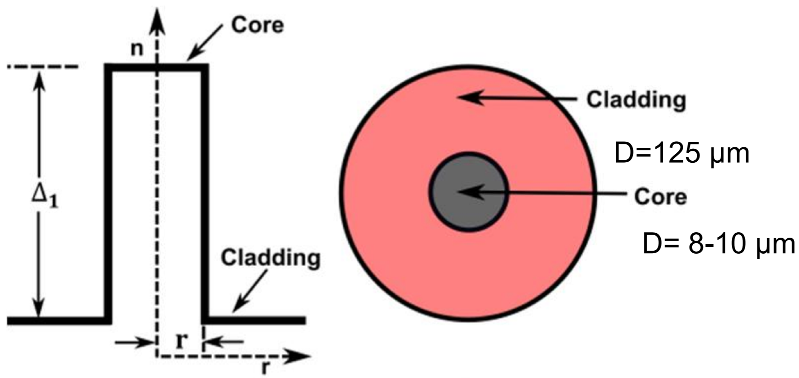
Time-Frequency metrology



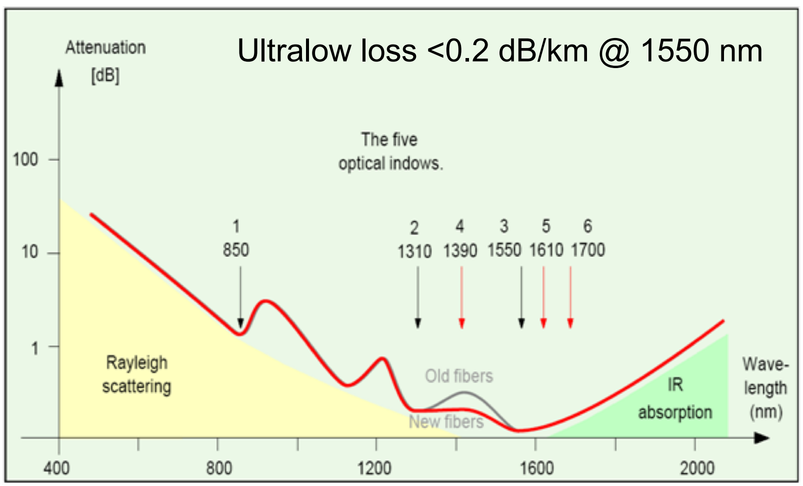
Imaging, Biophotonics

Basics of Fiber Optics

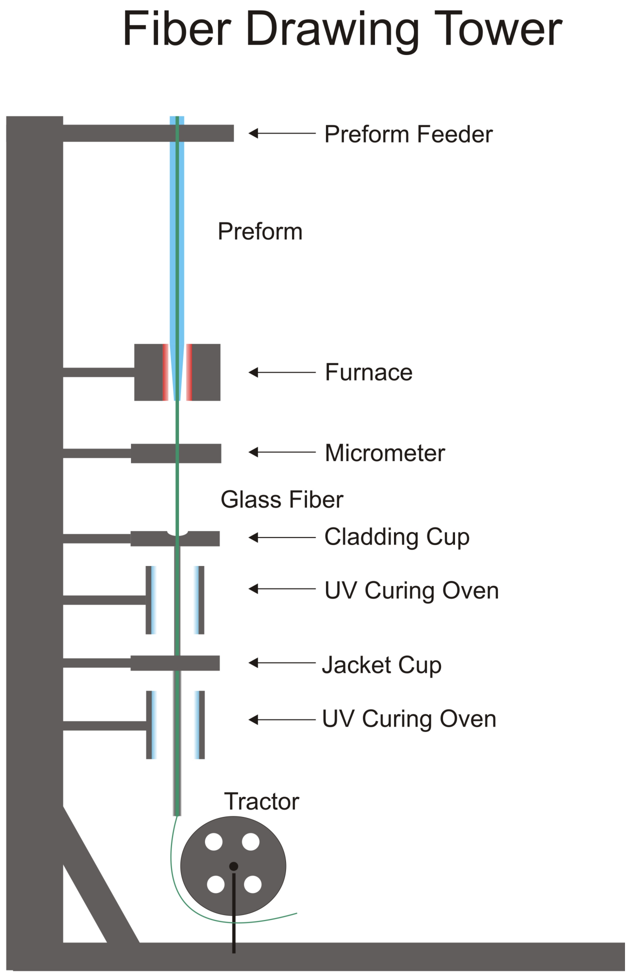
Typical refractive index profile and cross-section of a step-index fused silica (SiO_2) fiber



Attenuation spectrum of silica fibers

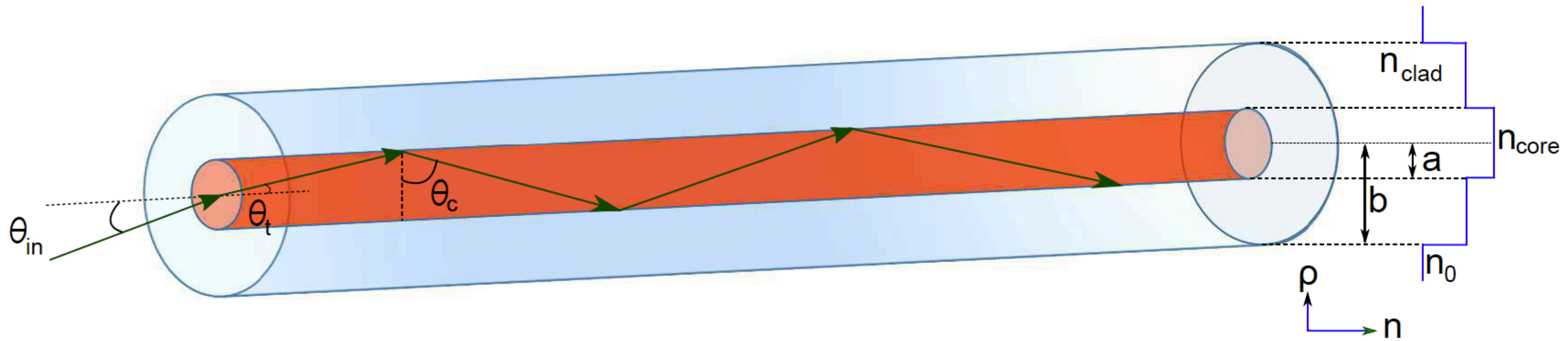


Manufacturing



- More than 2 billion kilometers installed fibers on our planet for communications !

Light guiding mechanism in standard step-index fibers by total internal reflection (TIR)

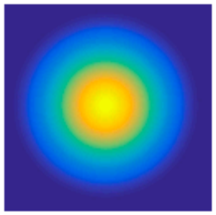


- Numerical Aperture (NA) $\sin(\theta_{\text{in}}) = \sqrt{n_{\text{core}}{}^2 - n_{\text{clad}}{}^2} = \text{NA}$
- Normalized frequency $V = p_c a = a k_0 (n_{\text{core}}{}^2 - n_{\text{clad}}{}^2)^{1/2} = \frac{2\pi a}{\lambda}\text{NA}$

✓ When *V is lower than V<2.405, the fiber can support only the fundamental mode LP01 which has the minimum loss of 0.2 dB/km at 1550 nm*

✓ If V>2.405, the waveguiding fiber is multimode

LP01 spatial mode

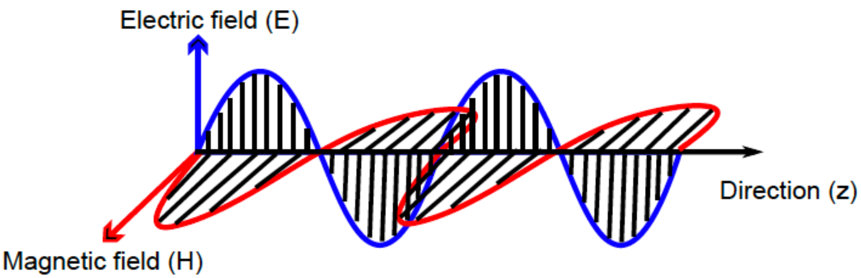


Basics of Fiber Optics

Light guiding is governed by the Helmholtz equations

- Electric field (FT)

$$\tilde{E}(\vec{r}, \omega - \omega_0) = F(x, y)\tilde{A}(z, \omega - \omega_0)\exp(i\beta_0 z).$$



$$\nabla_\perp^2 E + (k^2 - \beta^2)E = 0$$
$$\nabla_\perp^2 H + (k^2 - \beta^2)H = 0$$

$$k^2 = \frac{\omega^2}{c^2} n^2 = k_o^2 n^2$$

Wavevector

Propagation constant

$$\beta = \frac{2\pi}{\lambda_0} n_{\text{eff}}$$

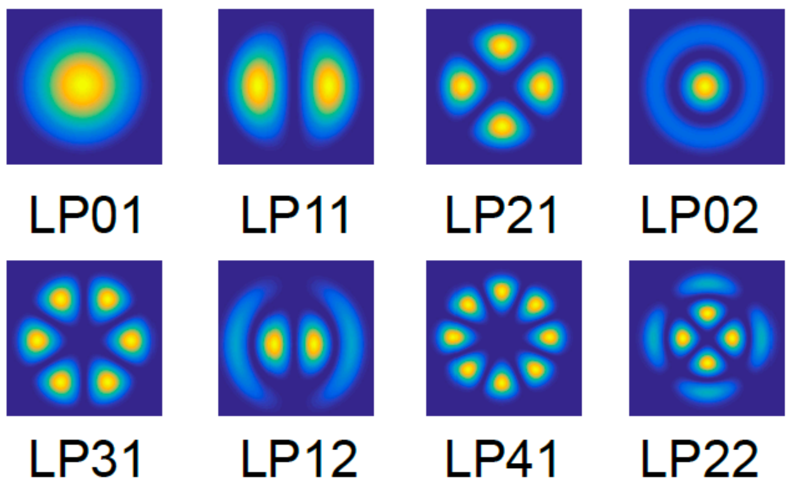
- Modal distribution (Bessel functions)

$$\frac{d^2F}{d\rho^2} + \frac{1}{\rho}\frac{dF}{d\rho} + \left(n^2 k_0^2 - \beta^2 - \frac{m^2}{\rho^2}\right)F = 0$$

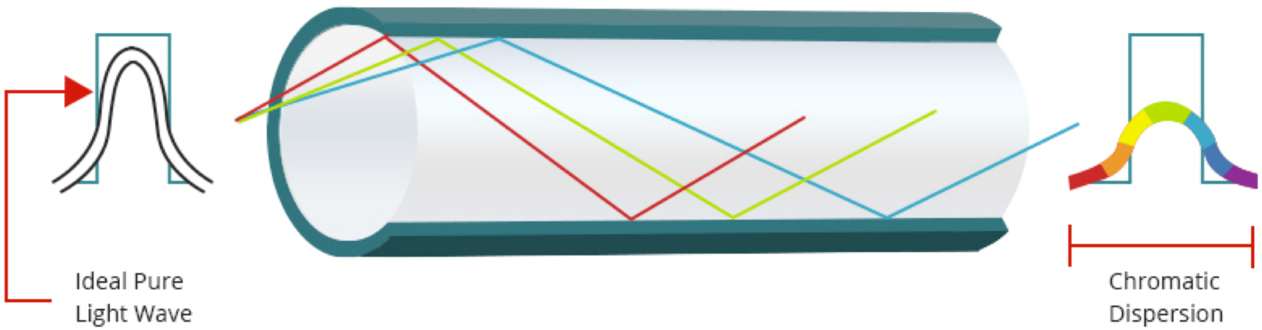
- Effective mode area (EMA)

$$A_{\text{eff}} = \frac{\left(\iint_{-\infty}^{\infty} |F(x, y)|^2 dxdy\right)^2}{\iint_{-\infty}^{\infty} |F(x, y)|^4 dxdy}$$

- Linearly-polarized (LP) fiber modes



Chromatic dispersion



- Signal or pulse spreading over time resulting from the different speeds of light rays
- A combination of the material and waveguide dispersion effects
- Material dispersion : wavelength dependence of the refractive index on the core material
- Waveguide dispersion : dependence of the propagation constant β on the fiber parameters
- Sellmeier equation for the refractive index

$$\boxed{n^2(\lambda) = 1 + \sum_{j=1}^{k} \frac{A_j \lambda^2}{\lambda^2 - B_j^2}}$$

$$\boxed{n_{\mathrm{eff}} = \frac{\beta}{k_0}}$$

Where A_j and B_j are the Sellmeier coefficients and n_{eff} the effective index

Chromatic dispersion

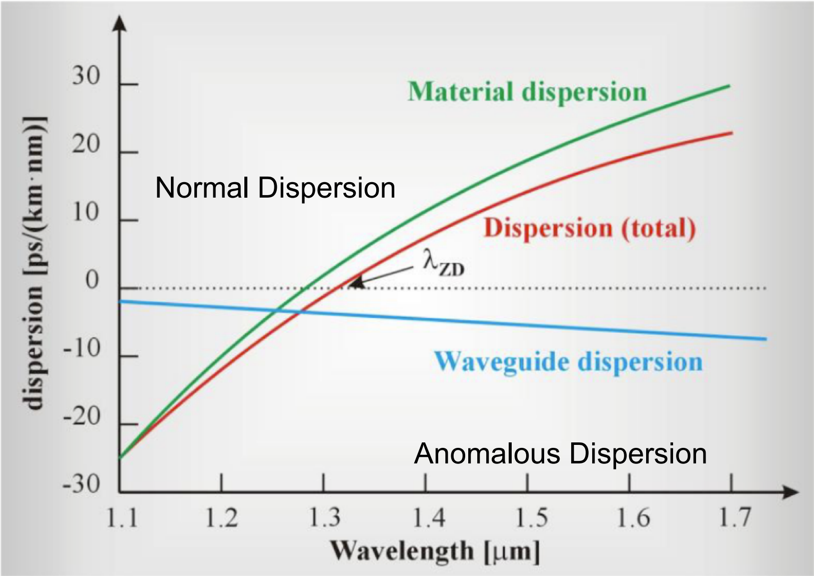
- Expanding the propagation constant as a Taylor series around the central frequency

$$\beta(\omega) = \beta_0 + \beta_1(\omega - \omega_0) + \frac{1}{2}\beta_2(\omega - \omega_0)^2 + \frac{1}{6}\beta_3(\omega - \omega_0)^3 + \ldots.. + \frac{1}{m!}\beta_m(\omega - \omega_0)^m$$

$$\beta_2 = \frac{1}{c}\left(2\frac{dn_{\text{eff}}}{d\omega} + \omega\frac{d^2n_{\text{eff}}}{d\omega^2}\right)$$ in s^2/m

$$D = \frac{d\beta_1}{d\lambda} = -\frac{2\pi c}{\lambda^2}\beta_2 = -\frac{\lambda}{c}\frac{d^2n_{\text{eff}}}{d\lambda^2}$$

D is in ps/nm/km

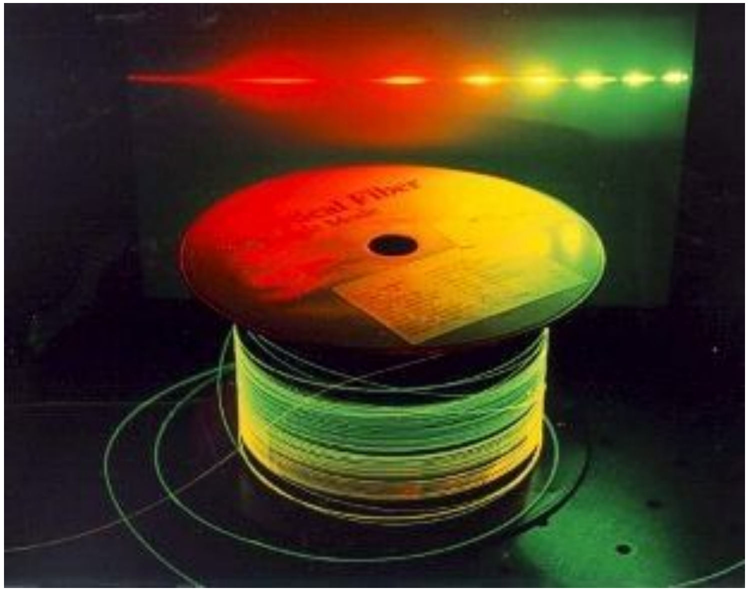


- When D is negative, D<0 => β_2>0: Normal dispersion regime (red light faster than blue)
- When D is positive, D>0 => β_2<0: Anomalous dispersion (blue light faster than red)
- When D is zero, D=0 => the zero-dispersion wavelength (ZDW)
- β_3 and β_4 are the dispersion slope and curvature

Raman Light Scattering

- Appearance of additional lines in the spectrum of monochromatic light that has been scattered by a transparent material medium, e.g. an optical fiber

Raman lines   Input Laser wavelength (532nm)

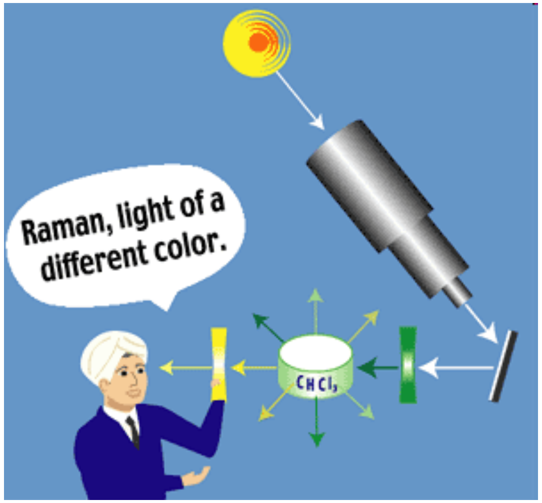


← Ouput spectrum

New colours !!!

Raman Light Scattering

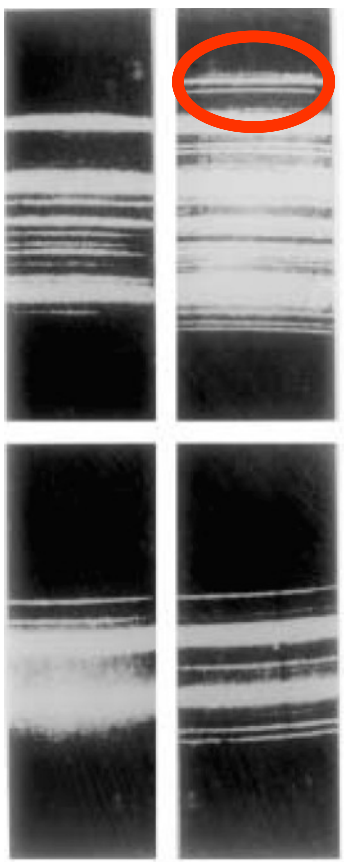
This effect was discovered by C. V. Raman in 1928 in **liquids** and simultaneously by G. Landsberg and L. Mandelstan in **crystals**.



C.V. Raman (1888 –1970) was an Indian physicist
Nobel Laureate recognised for his work on
the molecular scattering of light and the discovery of
« A new type of secondary radiation ». ***Nature*** (1928).
C. V. Raman remains the only Indian to receive a Nobel Prize in science.

Raman Light Scattering

Raman's article in *Nature* did not display any spectra.
They first appeared later in the *Indian Journal of Physics*.
Raman sent reprints of that article reporting his discovery to 2000 scientists in France, Germany, Russia, Canada, and the United States.



The first spectra taken by C. V. Raman and K. S. Krishnan

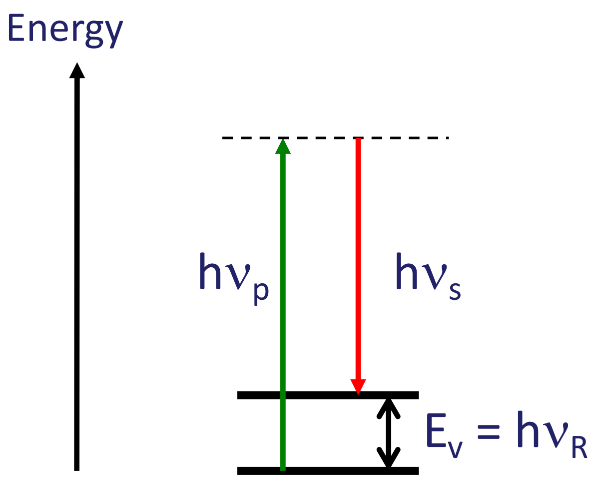
The upper-left photograph shows the incident light consisting of the spectrum of a quartz mercury arc lamp after passing through a blue filter that cuts out all wavelengths greater than the indigo line at 435.8 nm.

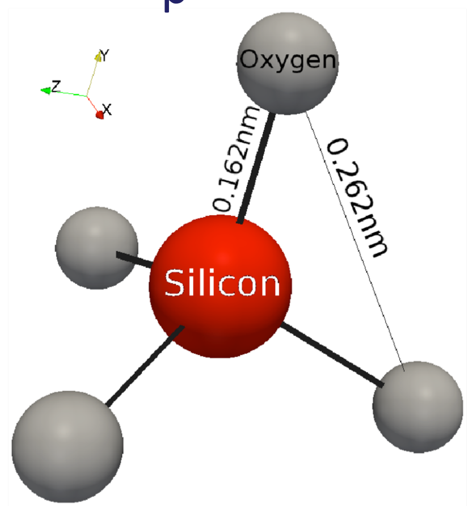
The upper-right photograph shows the same spectrum when scattered by liquid benzene and taken with a small spectroscope.
Note the appearance of modified lines owing to the Raman effect.

Spontaneous Raman scattering

Incoming photon

$$E_p = h\nu_p$$

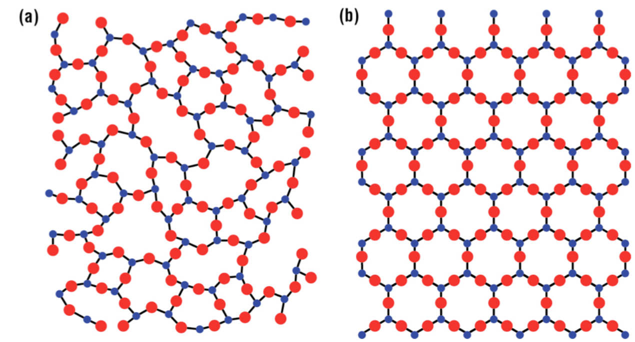




Silica Molecule

Silica Molecule (SiO_2)

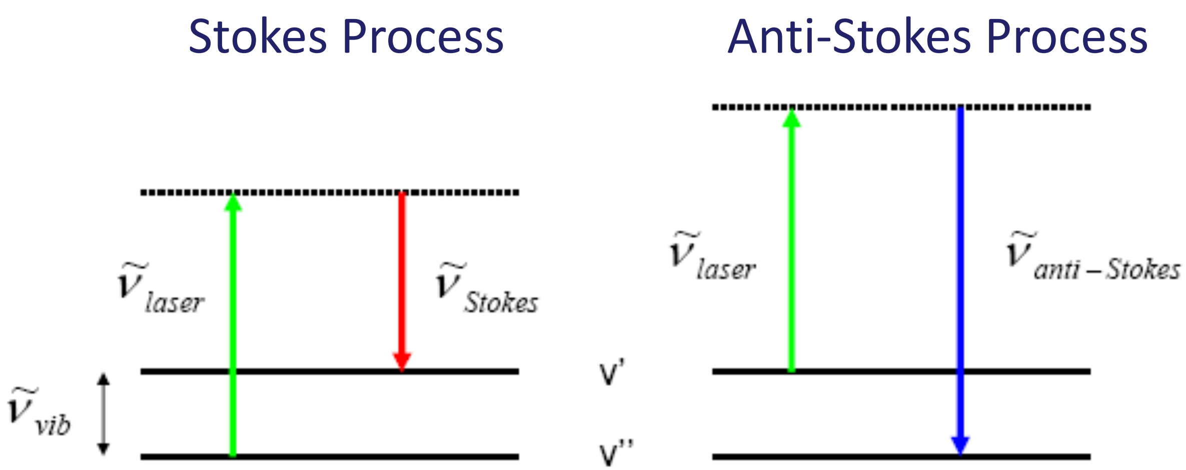
Amorphous Crystalline



$$E_p' = h\nu_p - E_v$$
$$= h\nu_s$$
$$= h(\nu_p - \nu_R)$$

Raman scattering is an ***inelastic scattering process***, where an incoming photon interacts with a coherently excited state of the system (e.g. the vibrational modes of a silica SiO_2 molecule). As a result of this interaction, a frequency down-converted (Stokes) photon is emitted.

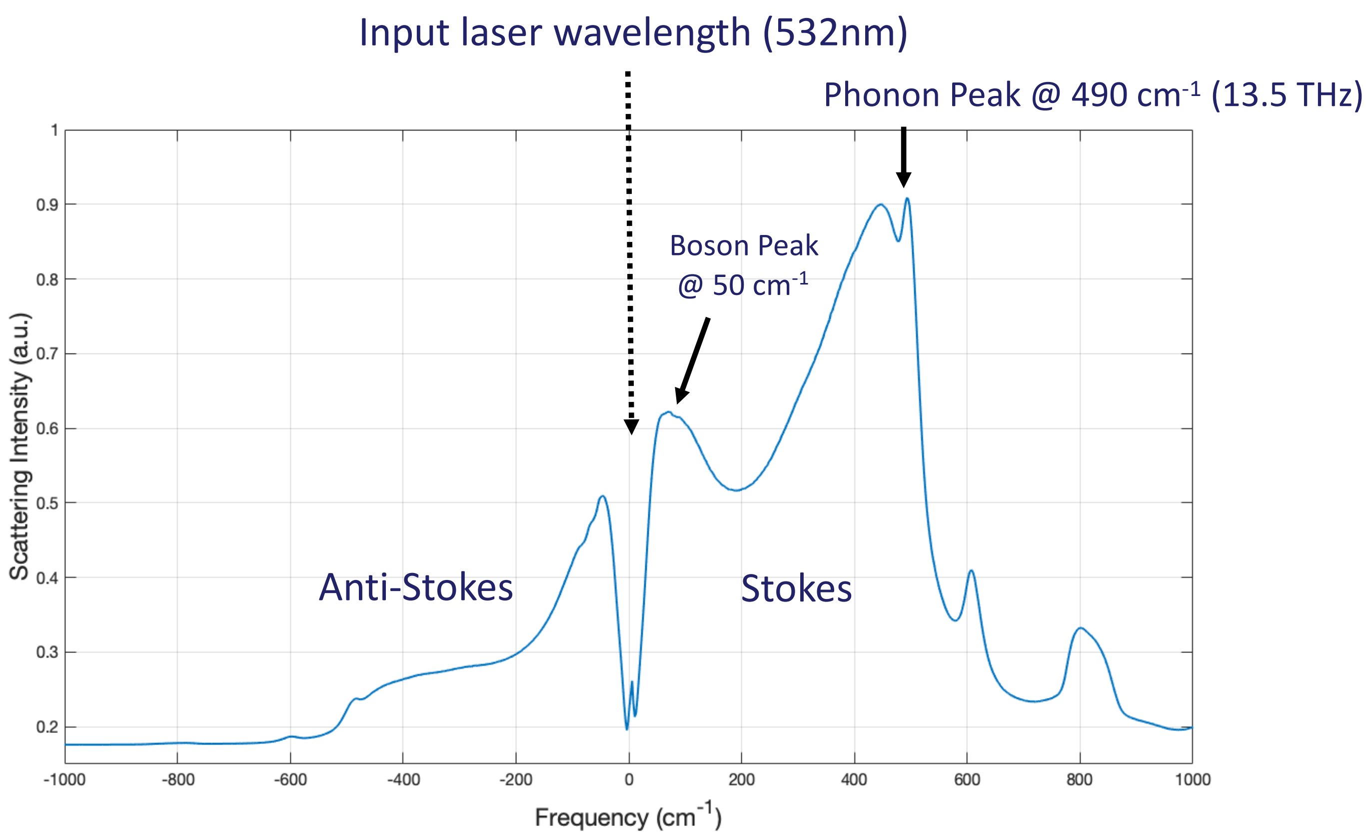
Spontaneous Raman scattering



$$\tilde{\nu}_{Stokes} = \tilde{\nu}_{laser} - \tilde{\nu}_{vib} \qquad \tilde{\nu}_{anti-Stokes} = \tilde{\nu}_{laser} + \tilde{\nu}_{vib}$$

An up-converted (anti-Stokes) photon may also be emitted if the vibrational states are sufficiently populated. However this is rarely observed in optical fibers due to the ultrafast relaxation time (fs).

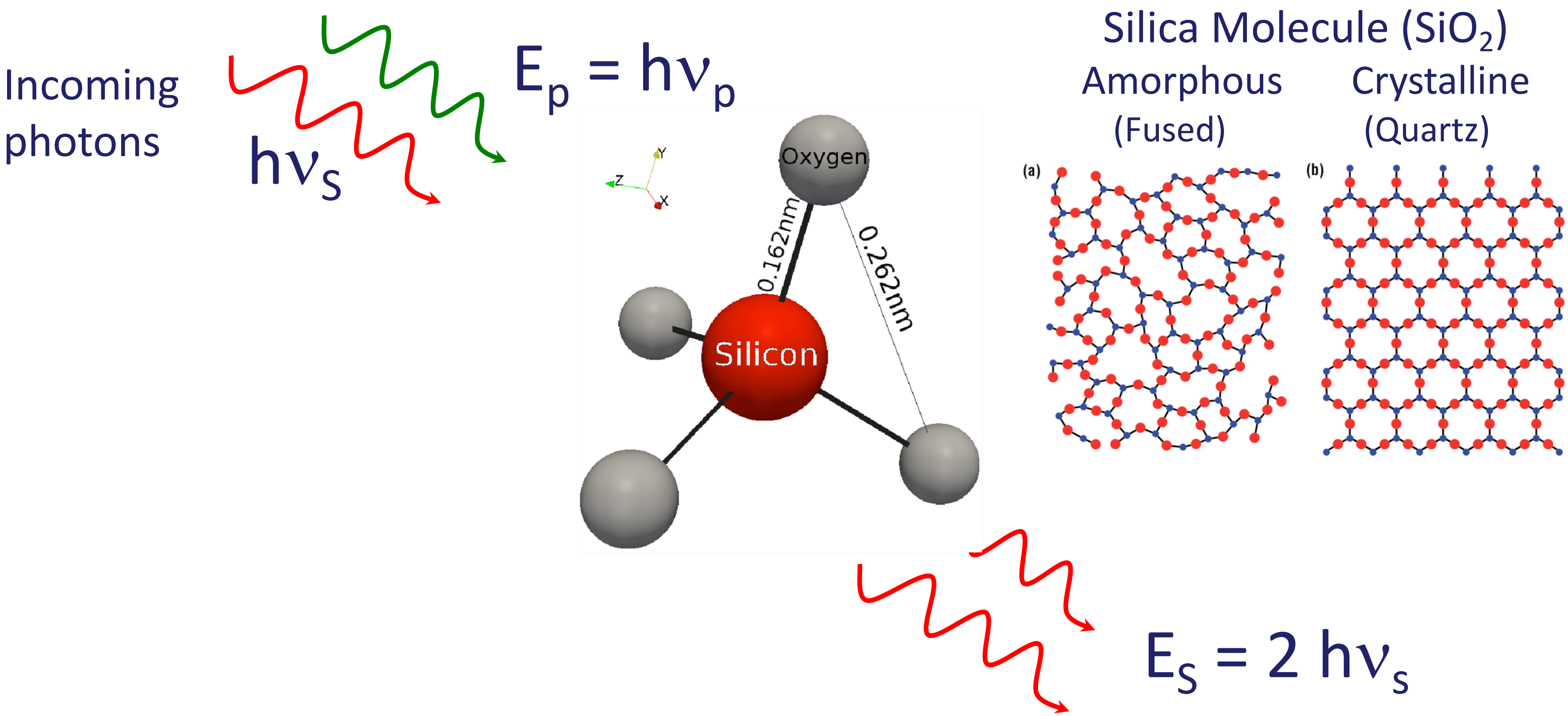
Spontaneous Raman scattering



The spontaneous Raman scattering efficiency is very weak and extends over a broad frequency range (1500 cm^{-1}). The ratio of the anti-Stokes to Stokes intensities follows the Maxwell-Boltzmann's distribution as:

$$I_S / I_{AS} \approx \exp\left(-\hbar\Omega_R / k_B T\right)$$

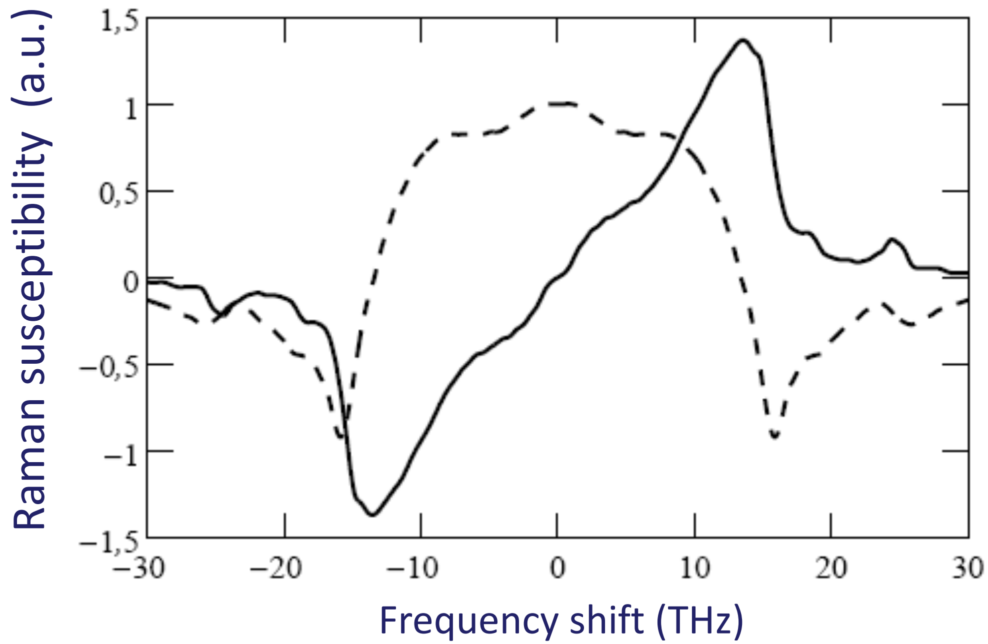
Stimulated Raman scattering (SRS)



If the incident laser beam is sufficiently intense, the photon-phonon scattering process becomes ***self-stimulated***. The temporal beating between the pump and the Stokes waves stimulates the vibrational states at the Raman frequency, and the wave grows rapidly such that most of ***the pump energy is transferred to it***.

Stimulated Raman scattering (SRS)

The Stimulated Raman gain spectrum in fused silica

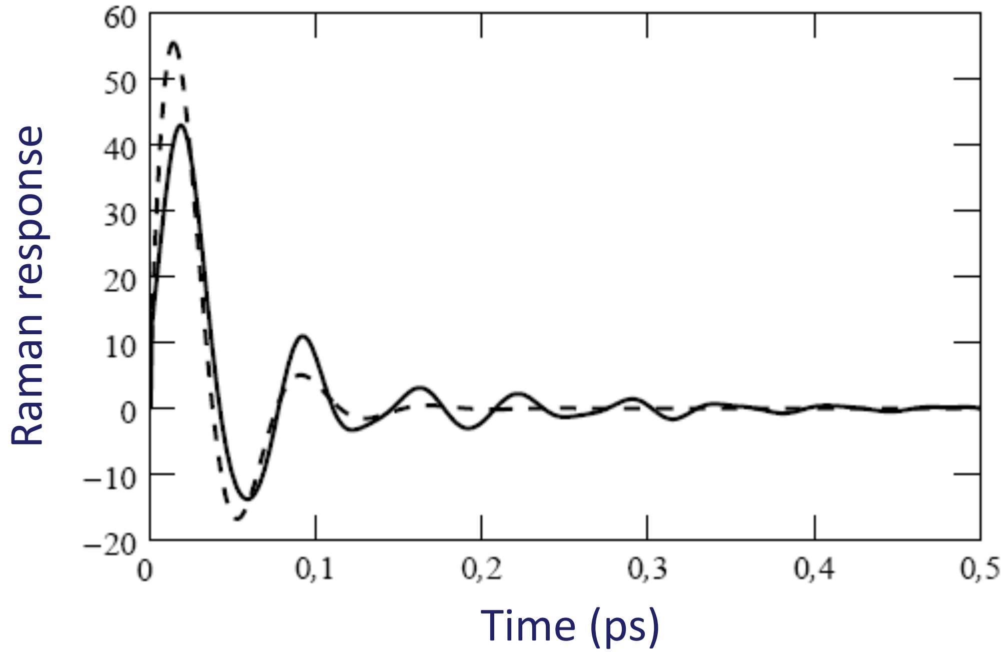


$$\chi_R^{(3)} = \chi_R^{'}(\omega) + i\chi_R^{''}(\omega) = FT(h_R(t))$$

The real part (symmetric dashed curve, nonlinear index) and imaginary part (antisymmetric solid curve, Raman gain) of the Raman susceptibility are related by the Kramers-Kronig relations.

Stimulated Raman scattering (SRS)

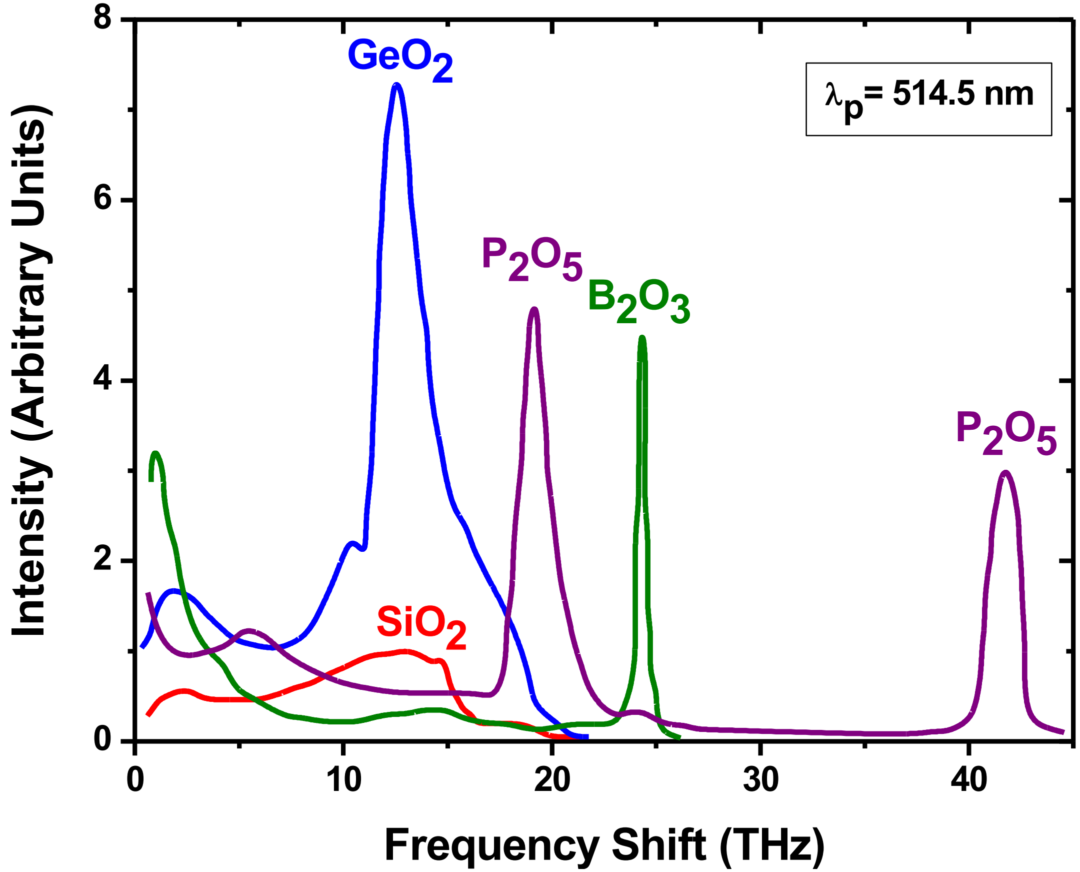
The Raman response function in the time domain is basically a decaying sinusoidal oscillation. The ***oscillation period*** corresponds to the ***peak of the Raman gain*** and the ***decay rate*** corresponds to the ***width*** of the gain spectrum.



$$h_R(t) = \frac{t_1^2 + t_2^2}{t_1 t_2^2} \sin\left(\frac{t}{t_1}\right) \exp\left(-\frac{t}{t_2}\right)$$

t_1=12.2 fs=$2\pi/\Omega_R$ (13.2 THz)
t_2=32 fs=$2\pi/\Delta\Omega_R$ (5 THz)

Stimulated Raman scattering (SRS)



Ref: Galeener et al., Appl. Phys. Lett. 32, 34 (1978)

The **Raman gain** can be **enhanced and frequency shifted** by use of dopants inside the silica core or using other nonlinear glasses (Chalcogenide, Tellurite, ZBLAN, fluoride)

Stimulated Raman scattering (SRS)

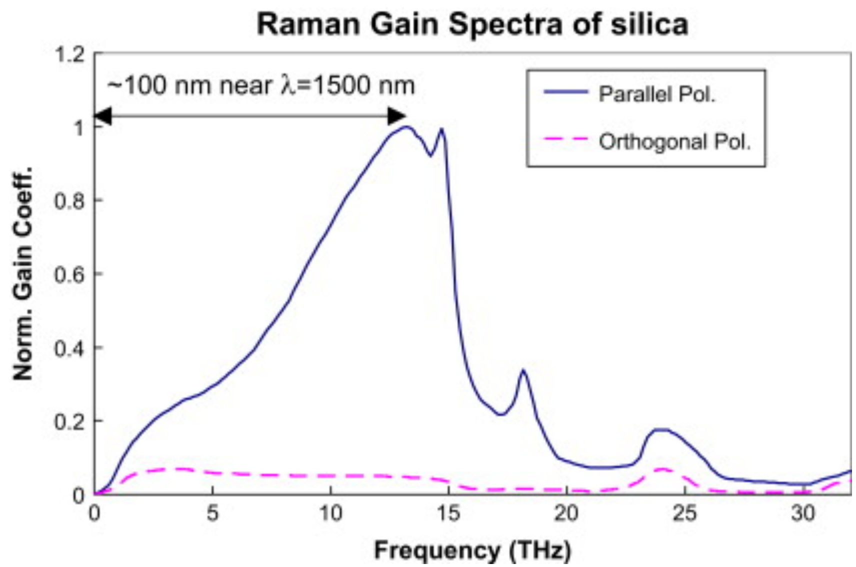
The Raman threshold is defined as the **input pump power** at which the **Stokes power** becomes equal to the **pump power** at the fiber output (Saturation regime) :

$$P_{Th}(W) = \frac{16 A_{eff}}{g_R L_{eff}}$$

- ✓ Where A_{eff} is the effective area of the HE11 mode : $A_{eff} = \frac{\left(\int\int_{-\infty}^{+\infty} |F(x,y)|^2 \, dx\, dy\right)^2}{\int\int_{-\infty}^{+\infty} |F(x,y)|^4 \, dx\, dy}$
- ✓ g_R is the Raman gain, typically 10^{-13}m.W^{-1} for silica fibers
- ✓ L_{eff} is the effective length that accounts for loss: $L_{eff} = \int_0^L e^{-\alpha z} \, dz = \frac{1 - e^{-\alpha L}}{\alpha}$

Stimulated Raman scattering (SRS)

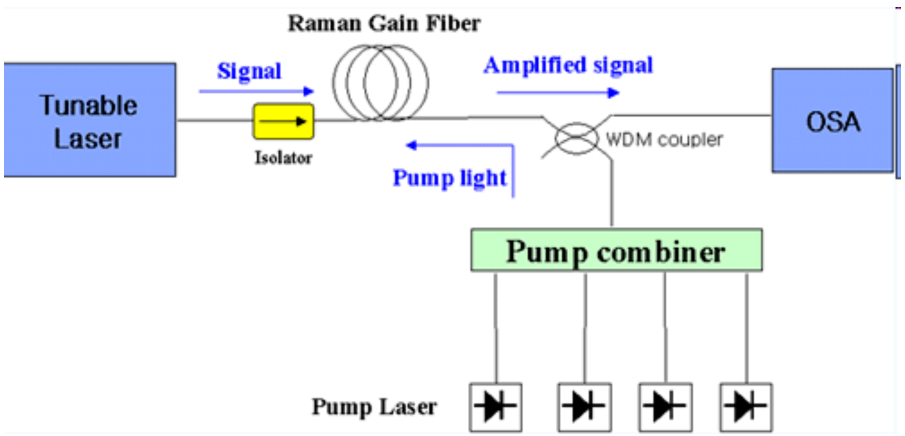
The effect of polarization state of pump light



Raman scattering is strongly polarization dependent. The Raman gain coefficient for a signal orthogonally polarized with respect to the pump is indeed ***30 times*** lower than the parallel Raman gain. However, by using cross-polarized pumps or pump polarization diversity, the Raman gain can be made independent of the state of polarization of the incident signal.

Key Applications : Fiber Raman Amplifiers

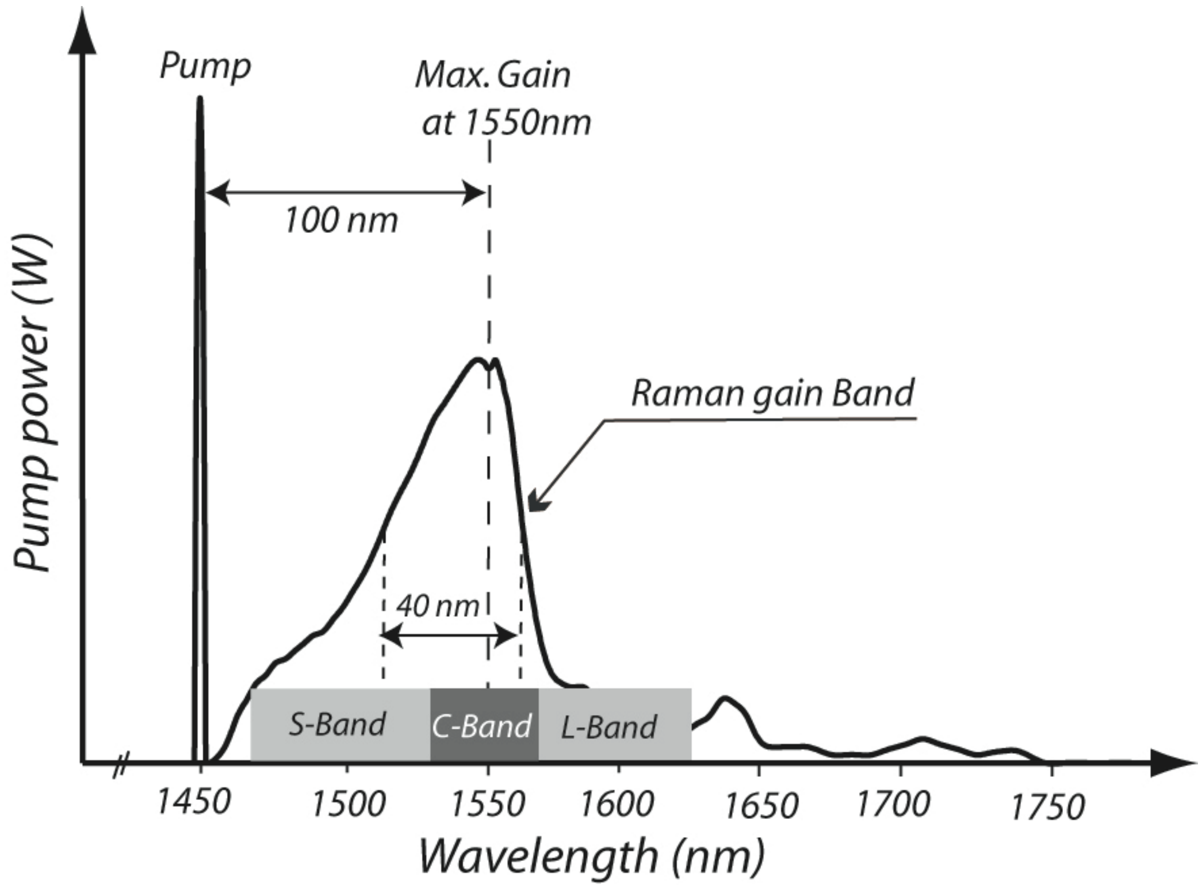
Stimulated Raman scattering can serve as an **optical amplifier** for applications in telecommunications systems (typically 50nm bandwidth @1.55µm).



Raman amplifiers can be operated in very different wavelength regions, provided that a suitable pump source is available. The gain spectrum can be tailored by using different pump wavelengths simultaneously. Distributed backward Raman fiber amplifiers can have a lower noise figure.

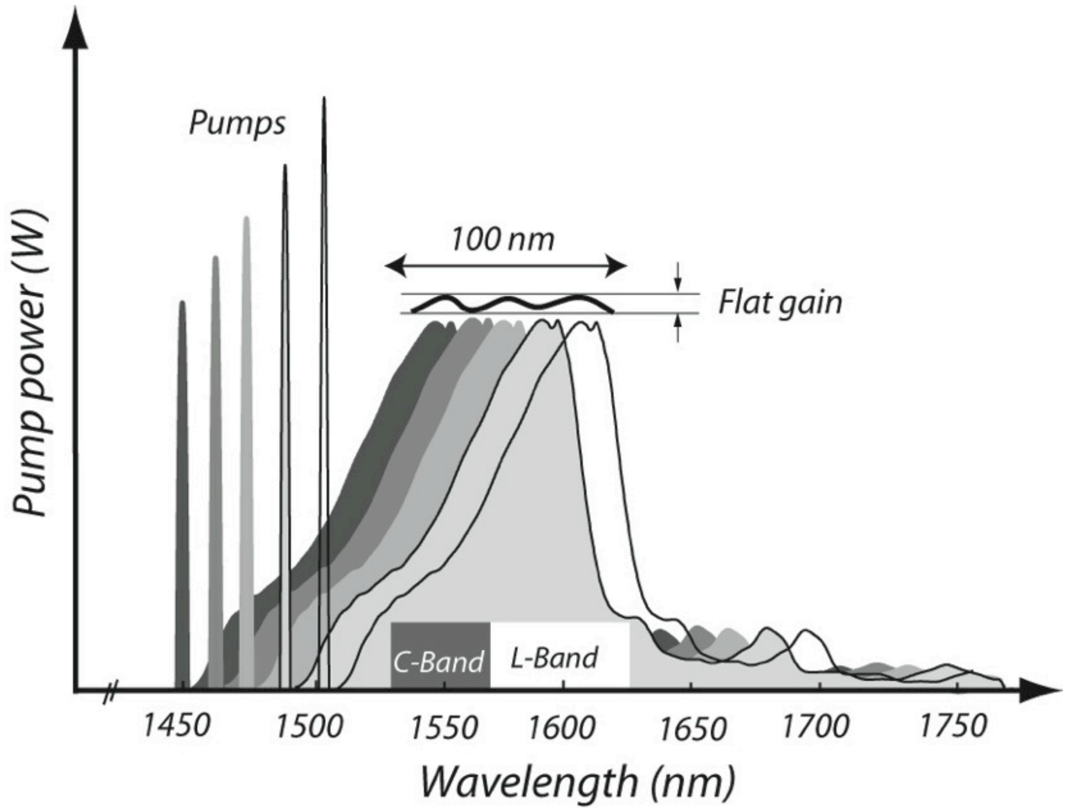
Key Applications : Fiber Raman Amplifiers

In optical fibers, the Raman effect is a broadband gain process.
This wide gain bandwidth is very advantageous for making Raman fiber amplifiers (RFA) and Raman fiber lasers (RFL). For a pump wavelength at 1450 nm, the gain band provides maximum gain at 1550 nm over a range of 40 nm.



Key Applications : Fiber Raman Amplifiers

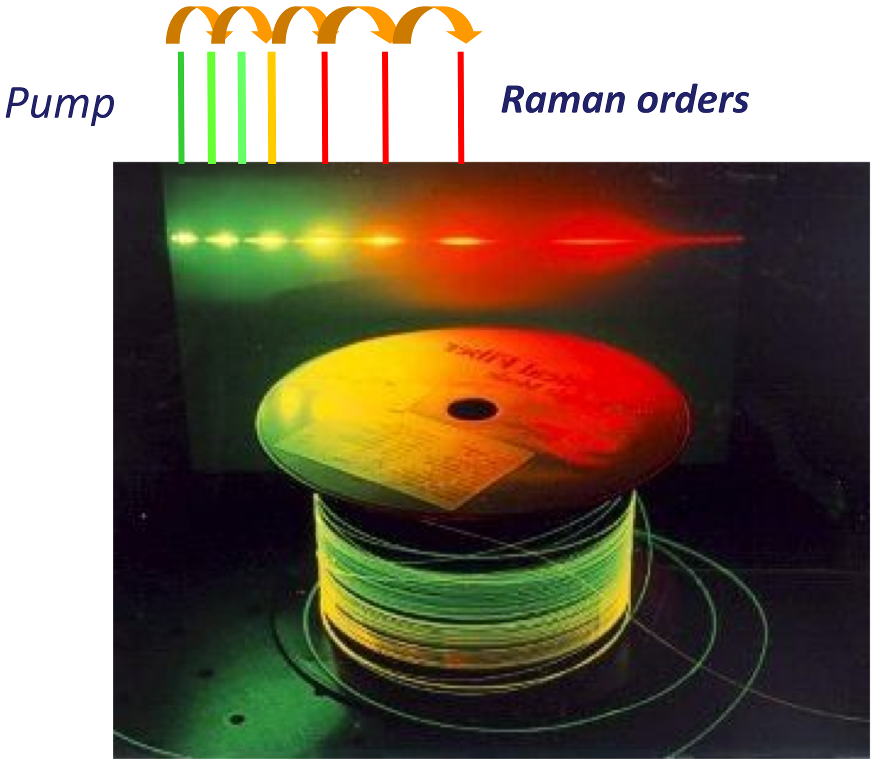
It is possible to combine in the same optical fiber several pump lasers such as to generate a flat and wide gain bandwidth. Such a large gain bandwidth has recently enabled to transmit more than 20 Tbit/s data rate over more than 10000 km.



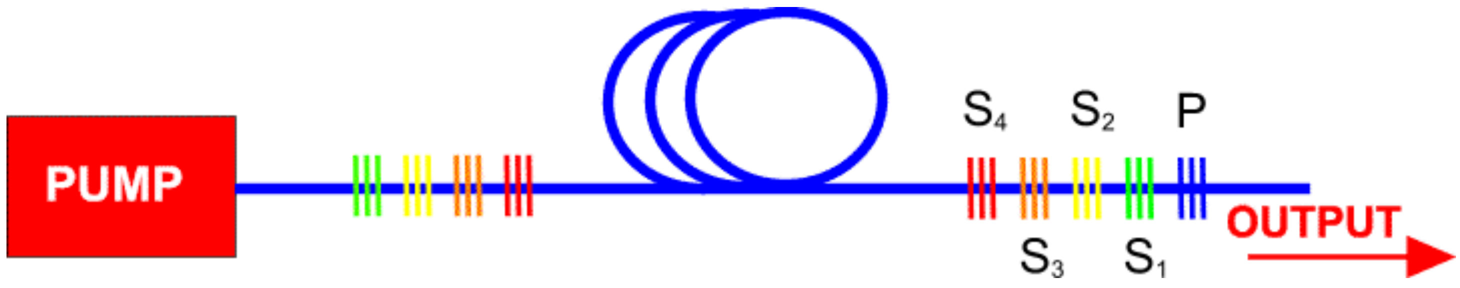
Principle of multi-wavelength pumping for flat and broadband Raman amplification at telecommunication wavelengths (C and L bands).

Cascaded Raman Scattering (CRS)

- Cascaded Raman generation is an iteration of fundamental stimulated Raman scattering (SRS) processes in which each generated Stokes wave acts as a pump to produce an **additionnal order**



Cascaded Raman fiber lasers (RFL)



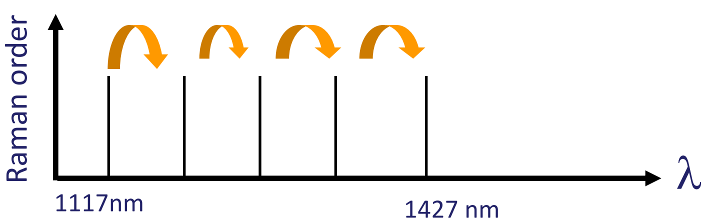
Cascaded Raman fiber lasers can be built with nested pairs of fiber Bragg gratings. Oscillation on one Raman order is used for pumping another order, so that larger frequency offsets can be bridged.

▸**Advantages**

☺ Very high output power delivery

☺ Wavelength versatile

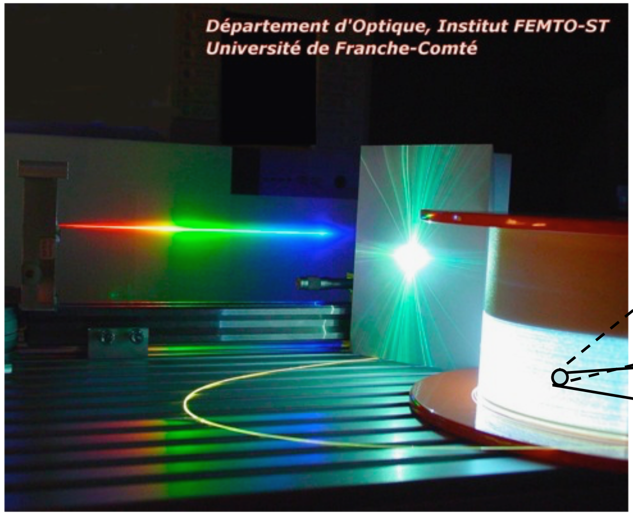
☺ Ideal for Raman amplification in WDM systems (low noise, wideband,...)

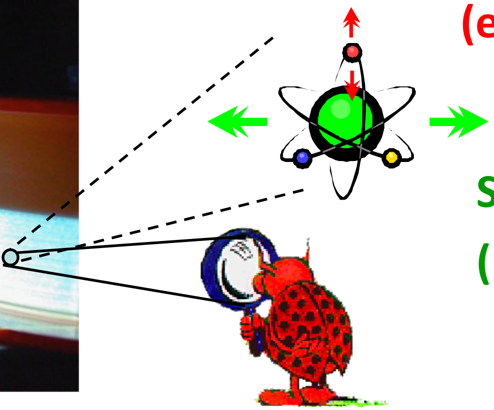


Principle of operation

The Optical Kerr Effect

Nonlinear optical effects are striking phenomena that arise when an intense optical beam propagates through an optical fiber





Optical Kerr Effect (electronic response)

$n = n_0 + n_2 I(z,t)$

Stimulated Raman Scattering (Delayed Molecular response)

Third-Order Nonlinear Polarization

$$P_{NL}(t) = \boxed{\varepsilon_0 \chi_K^{(3)} : E(t)E(t)E(t)} + \boxed{\varepsilon_0 E(t) \int_{-\infty}^{t} \chi_R^{(3)}(t-t')E(t')E(t')dt'}$$

Optical Kerr effect:
Instantaneous Elastic effect
(No energy exchange with matter)

Raman effect :
Inelastic light scattering:
(Molecular vibrational state)

The Optical Kerr Effect

The optical **Kerr effect** can be described as an instantaneous and local change in the refractive index, proportional to the optical intensity:

$$\mathbf{n(z,t)=n_0+n_2I(z,t)}$$


John Kerr
(1824-1907)

with **$\mathbf{n_2}$ the nonlinear Kerr index**

(typically $3 \times 10^{-16} cm^2/W$ for silica)

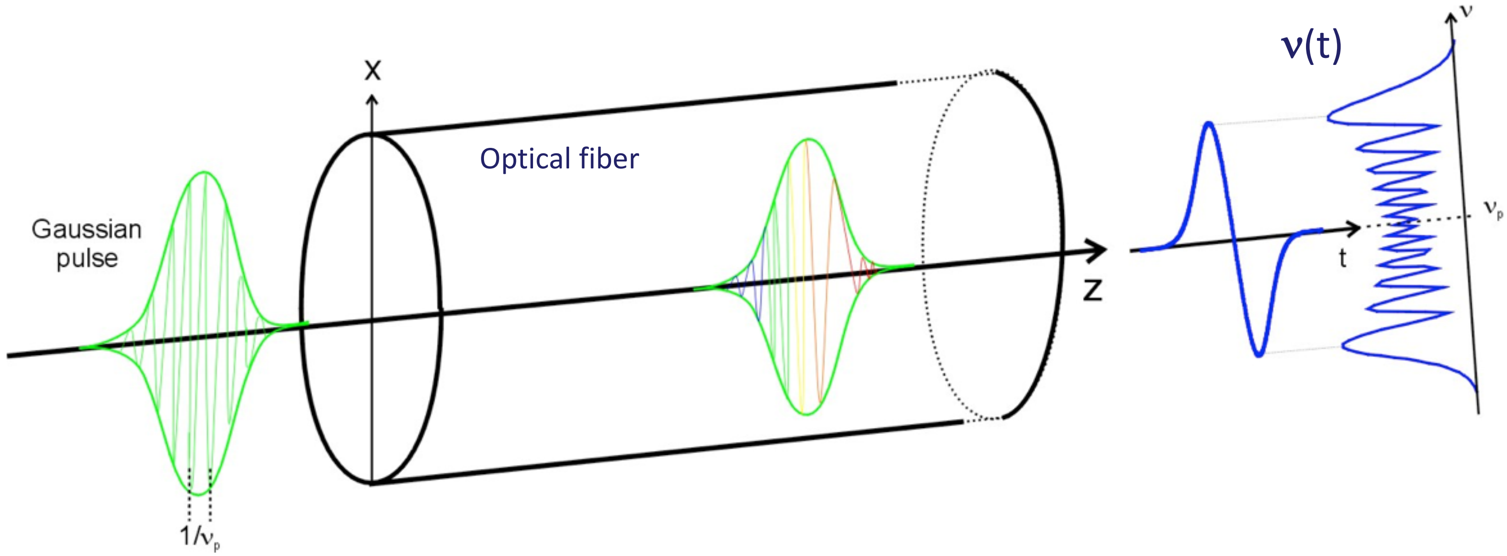
In optical fibers, we use the nonlinear coefficient : $\gamma = \dfrac{2\pi n_2}{\lambda A_{eff}}$ ($W^{-1}\,km^{-1}$)

This is an elastic (non-resonant) scattering : the incident photons suffer phase and/or frequency shifts **but overall there is no energy exchange with the material**

- Kerr materials exhibit many quantum properties: noise squeezing, correlation, twin-photon, entanglement, EPR paradox, etc…

Self-phase modulation (SPM)

If an optical pulse is transmitted through an optical fiber, the Kerr effect causes a time-dependent phase shift according to the pulse profile. In this way, the pulse acquires a so-called chirp, i.e., a temporally varying instantaneous frequency, and therefore a spectral broadening.

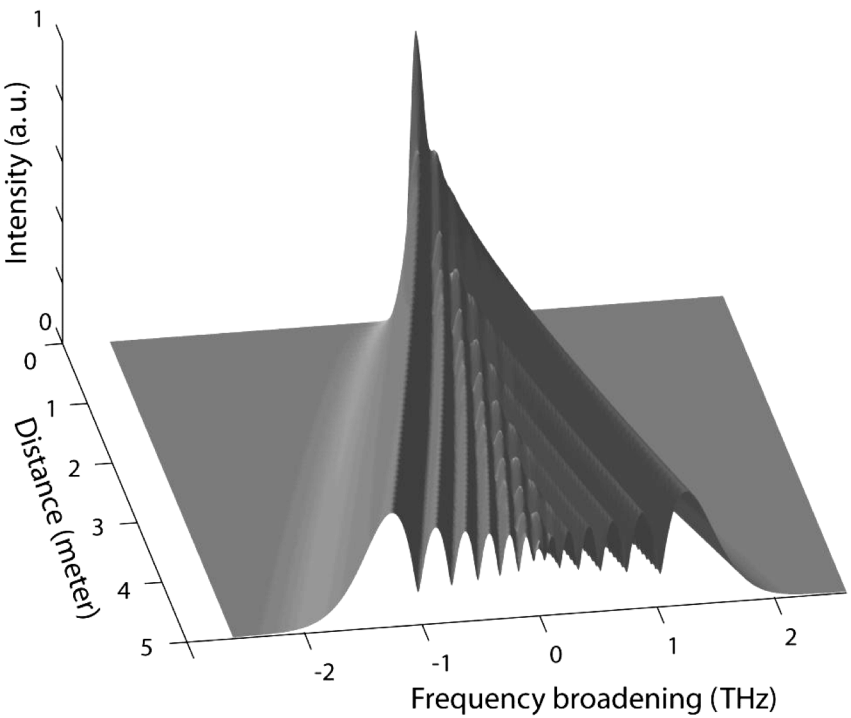


Nonlinear phase shift: $\varphi_{NL}(t) = \gamma P(t) L$

Instantaneous frequency chirp: $\nu(t) = -\dfrac{1}{2\pi}\dfrac{\partial \varphi_{NL}}{\partial t} = -\dfrac{\gamma L}{2\pi}\dfrac{\partial P(t)}{\partial t}$

Self-phase modulation (SPM)

The Fourier spectrum of the optical pulses exhibits strong oscillations due to constructive and destructive interferences



$$A(z,t) = A(0,\tau)\exp\left(j\gamma L|A|^2\right)$$

$$\left|A(\tau)\right|^2 = P_0 \exp\left(-\tau^2/\tau_0^2\right)$$

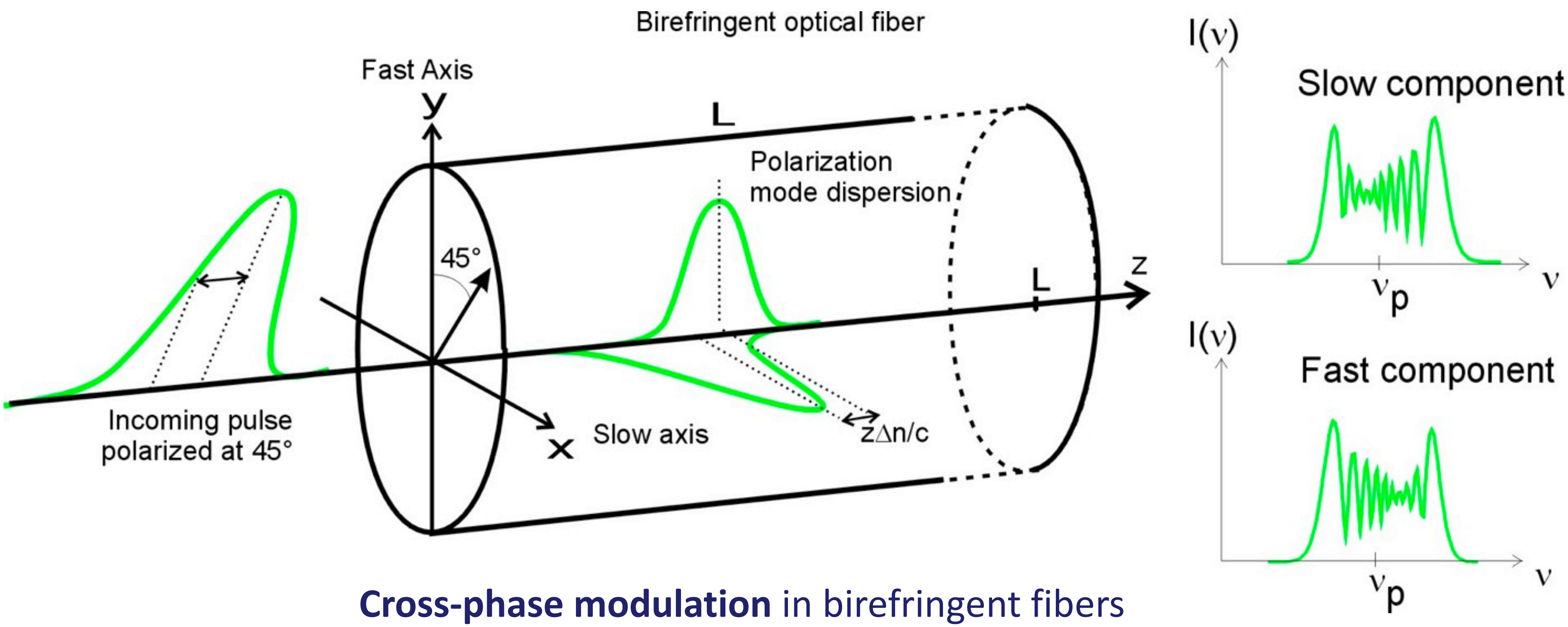
$$\delta\omega(t) = -\gamma L \frac{\partial}{\partial t}|A|^2 = 2\gamma P_0 L \tau \exp\left(-\frac{\tau^2}{\tau_0^2}\right)$$

The number of fringes *N* is proportional to the nonlinear phase shift by $N\pi$.

R.H. Stolen and C.H. Lin, Phys. Rev. A 17, 1448-1453 (1978).

Cross-phase modulation (XPM)

Cross-phase modulation (XPM) is the change in the optical phase of a light pulse caused by the interaction with another pulse of different color or polarization. Compared with self-phase modulation, there is an additional factor of 2 and for cross-polarized beams in birefringent fiber, it must be replaced with 2/3.



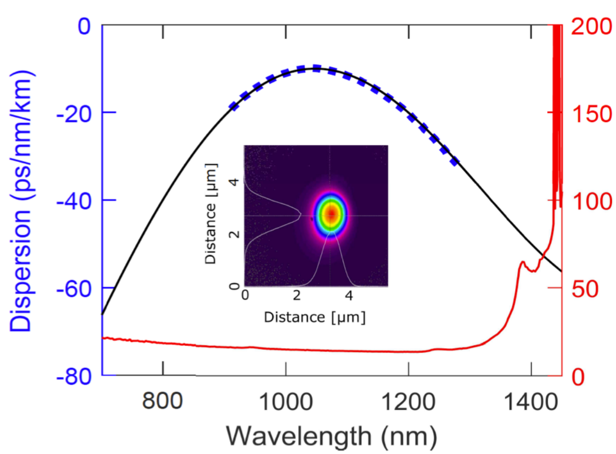
Cross-phase modulation in birefringent fibers

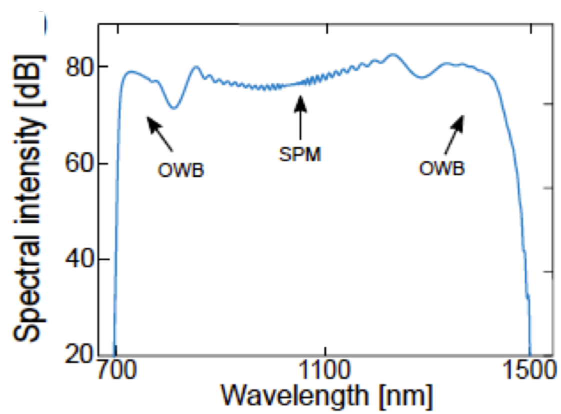
Optical Wave Breaking (OWB)

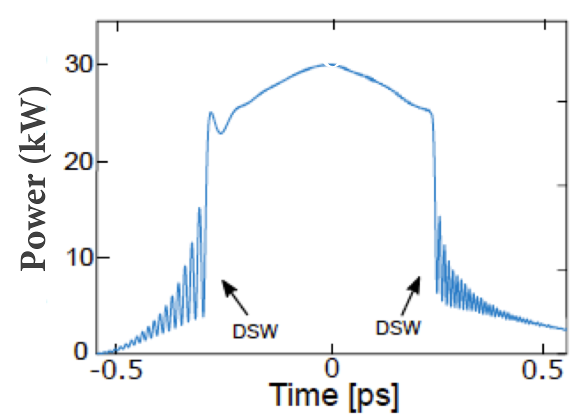
- Modelling with a generalized nonlinear Schrödinger equation (NLSE)

$$\underbrace{\frac{\partial A}{\partial z} = -\frac{\alpha(\omega)}{2}A}_{\textit{loss}} + \underbrace{\sum_{k\geq 2}\frac{i^{k+1}}{k!}\beta_k\frac{\partial^k A}{\partial T^k}}_{\textit{dispersion}} + i\gamma\underbrace{\left(1 + i\tau_0\frac{\partial}{\partial T}\right)}_{\textit{Self-steepening}}\underbrace{\left(A\int_{-\infty}^{+\infty} R(T')|A(z,T-T')|^2 dT'\right)}_{\textit{Nonlinear response (Kerr + Raman)}}$$

- Pumping in the normal dispersion regime
- Self-phase modulation (SPM) + Optical-Wave Breaking (OWB)
- OWB = Dispersive Shock Wave (DSW) => FWM between SPM and residual pump
- Smooth and flat SC spectrum – Pulse-preserved
- Gives low noise and high coherence level







A Heidt et al (2017). JOSAB, 34(4), 764. *C. Finot et al, JOSAB 25, 1938 (2008).*

Thibaut.sylvestre@univ-fcomte.fr

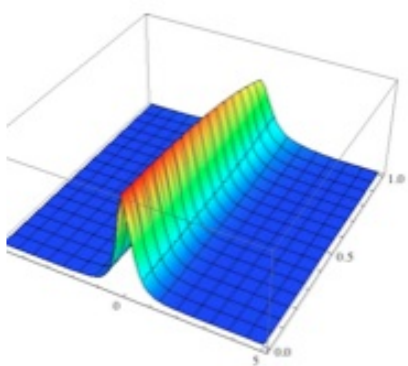
Temporal Optical Soliton

A fundamental optical soliton (N=1) is a pulse that preserves its shape during the propagation, being unaffected by dispersion and nonlinearities or collision.

Soliton amplitude: $A(z,t) = \sqrt{P_0}\,\mathrm{sec}h\left(\frac{\tau}{\sqrt{\beta_2/\gamma P_0}}\right)\exp\left(i\frac{\gamma P_0}{2}z\right)$

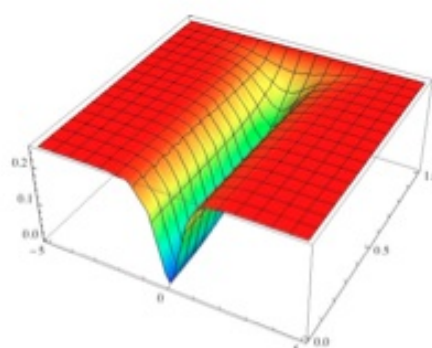
Soliton number: $N = \sqrt{\frac{\gamma P_0 T_0^2}{|\beta_2|}} = 1$

Bright Soliton



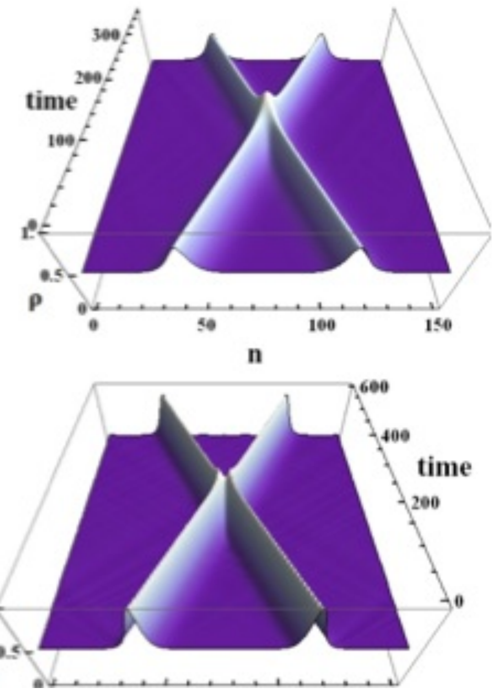
Anomalous dispersion (β_2<0)

Dark Soliton



Normal dispersion (β_2>0)

Elastic collision

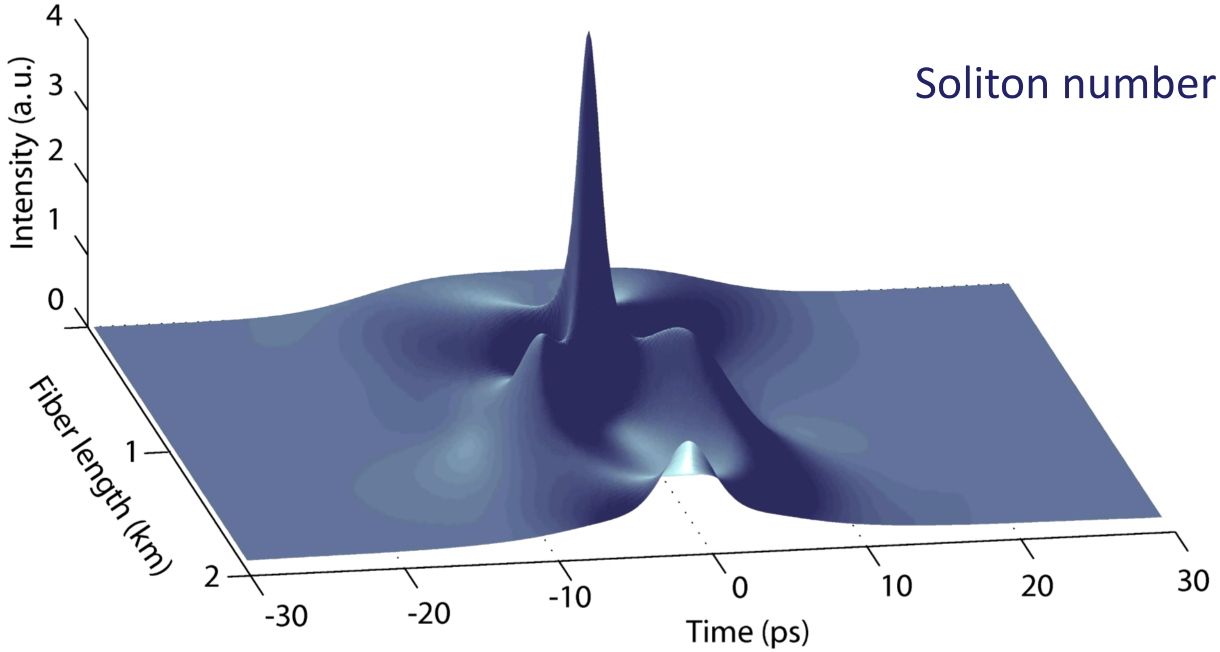


Solitons cross without interaction

High-order Optical Soliton

A high-order optical soliton (N>1) is a pulse that periodically changes its shape during the propagation, always returning to its original shape.

Soliton number: $N = \sqrt{\dfrac{\gamma P_0 T_0^2}{|\beta_2|}} = 3$



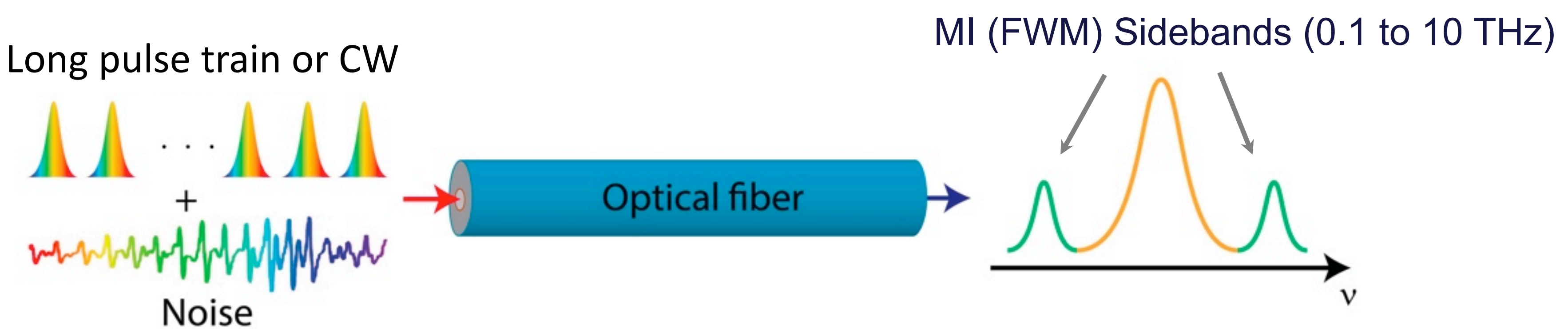
Third-order soliton (N=3) propagating in a 2 km-long single-mode fiber

Modulation Instability (MI)

Modulation instability manifests as the break-up of a continuous field or long pulses into a train of ultra-short pulses (fs).

MI gain: $g(\Omega) = |\beta_2 \Omega| \sqrt{\frac{4\gamma P_0}{|\beta_2|} - \Omega^2}$

MI modulation frequency : $\Omega_{\max} = \pm \sqrt{\frac{2\gamma P_0}{|\beta_2|}}$

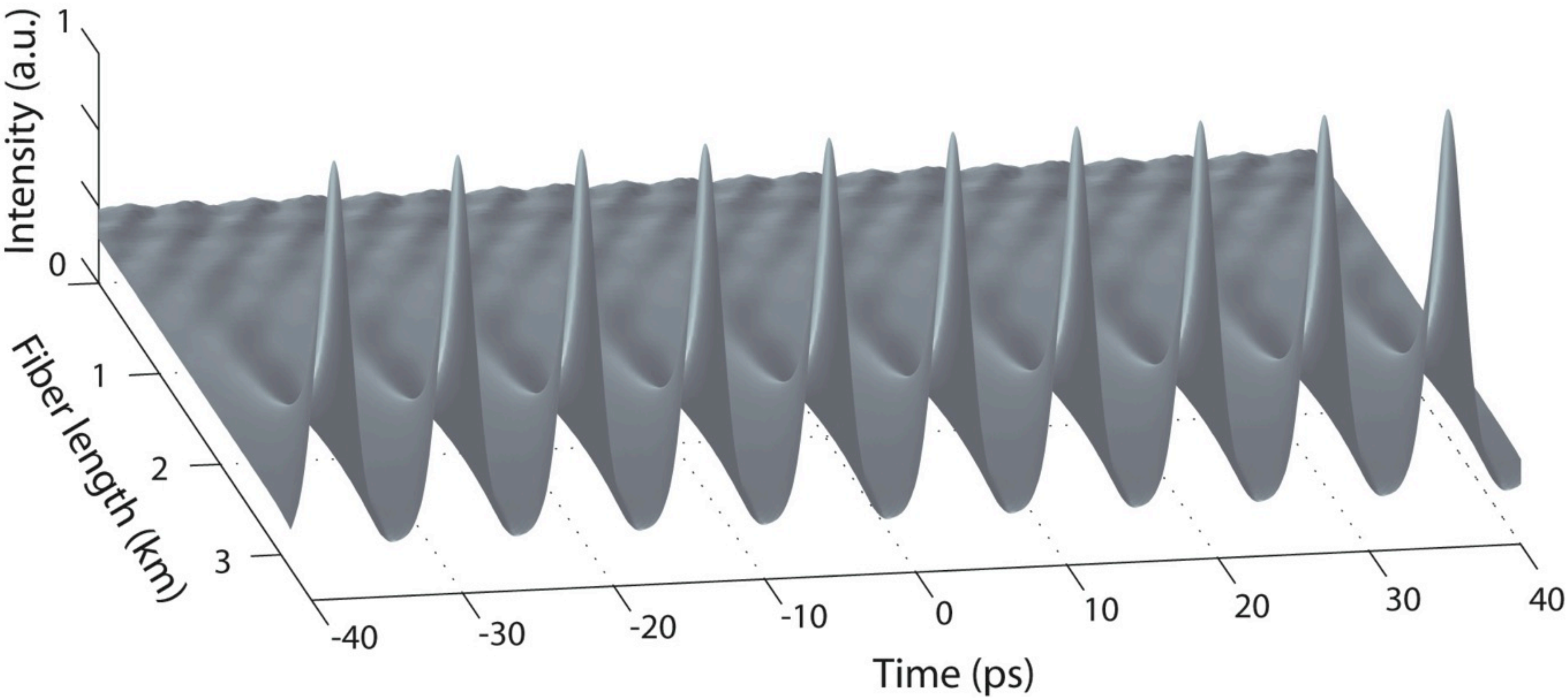


In the frequency domain, it gives rise to two symmetric FWM sidebands (Stokes and anti-Stokes).

Modulation instability plays a fundamental role in supercontinuum generation with cw or long pulse pumping.

Modulation Instability (MI)

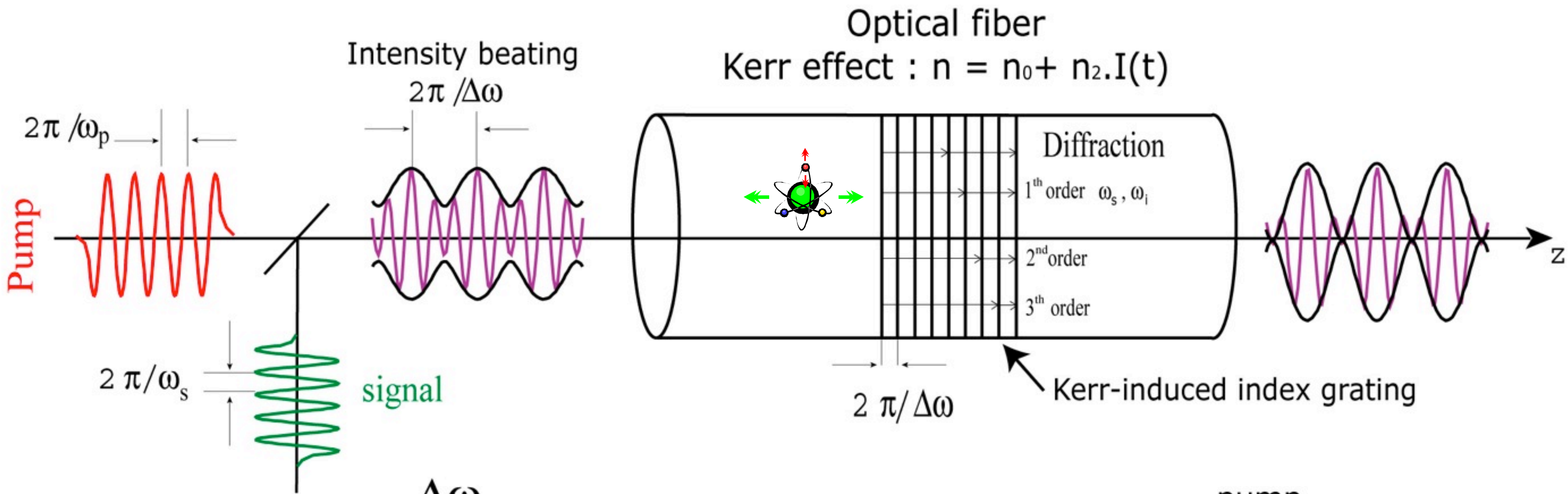
Modulation instability manifests as the break-up of a continuous field or long pulses into a train of ultra-short pulses (fs).

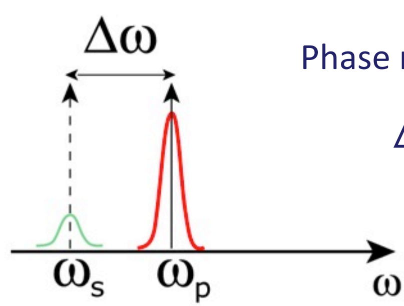


Modulation instability of a noisy continuous field and a soliton-like (breathers) pulse train generated in a 3-km optical fiber, obtained from a numerical simulation of NLSE

Optical Parametric Amplification (OPA)

A phenomenological approach based on interference and diffraction

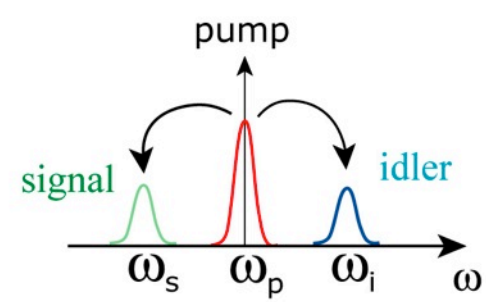




Phase matching ⇔ grating step

$$\Delta\beta = \Delta\beta_L + \Delta\beta_{NL} = 0$$

$$2\ \omega_p = \omega_s + \omega_i$$



Parametric gain: $G = 1 + \left(\frac{\gamma P}{g} \sinh\left(g L_{eff} \right) \right)^2$ with $g^2 = \left(\gamma P\right)^2 - \left(\kappa/2\right)^2$

Parametric amplification ⇔ Four-wave mixing ⇔ Modulation Instability ⇔ optical solitons

Optical Parametric Amplification (OPA)

Phase-matching condition near the zero-dispersion wavelength

$$\kappa = \Delta\beta_L + \Delta\beta_{NL}$$

$$\Delta\beta_L = \beta_2 \Delta\omega^2 + \frac{\beta_4}{12}\Delta\omega^4$$

Linear phase shift

$$\Delta\beta_{NL} = 2\gamma P_0 > 0$$

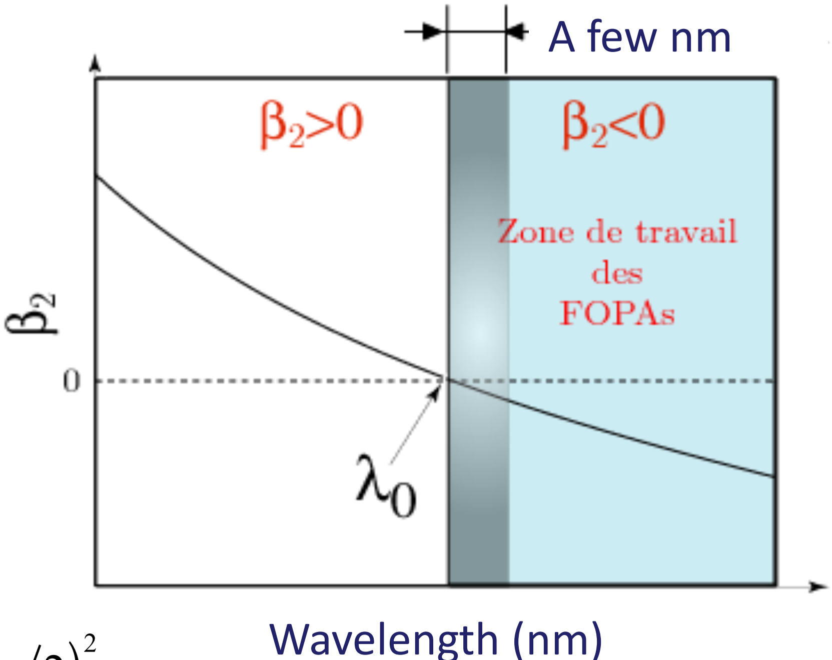
Nonlinear phase shift

⇨ **Dispersion and nonlinearity must cancel out**

⇨ **Amplification if β_2 and/or $\beta_4 < 0$**

- **if λ_P close to λ_0 (1.55 μm)**

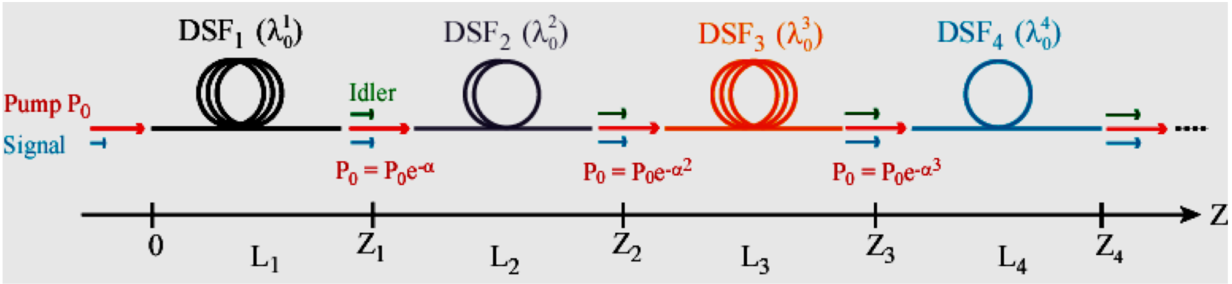
⇨ **Achievement of ultra-wide gain bandwidth**

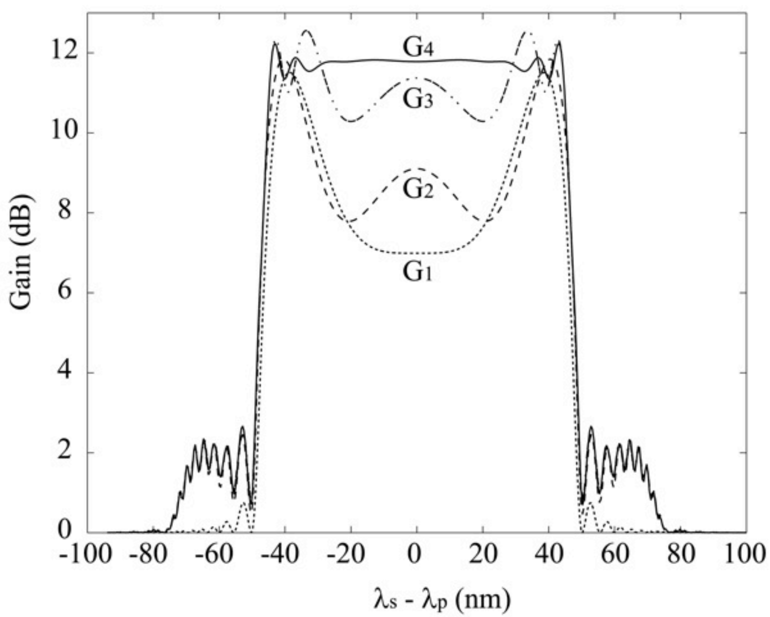


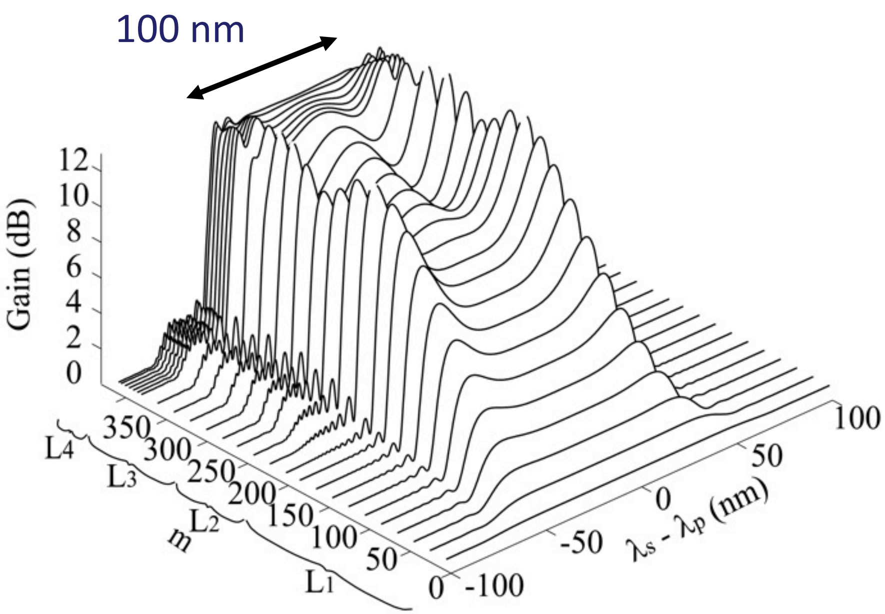
$$G = 1 + \left(\frac{\gamma P}{g}\sinh\left(gL_{eff}\right)\right)^2 \quad g^2 = \left(\gamma P\right)^2 - \left(\kappa/2\right)^2$$

Optical Parametric Amplification (OPA)

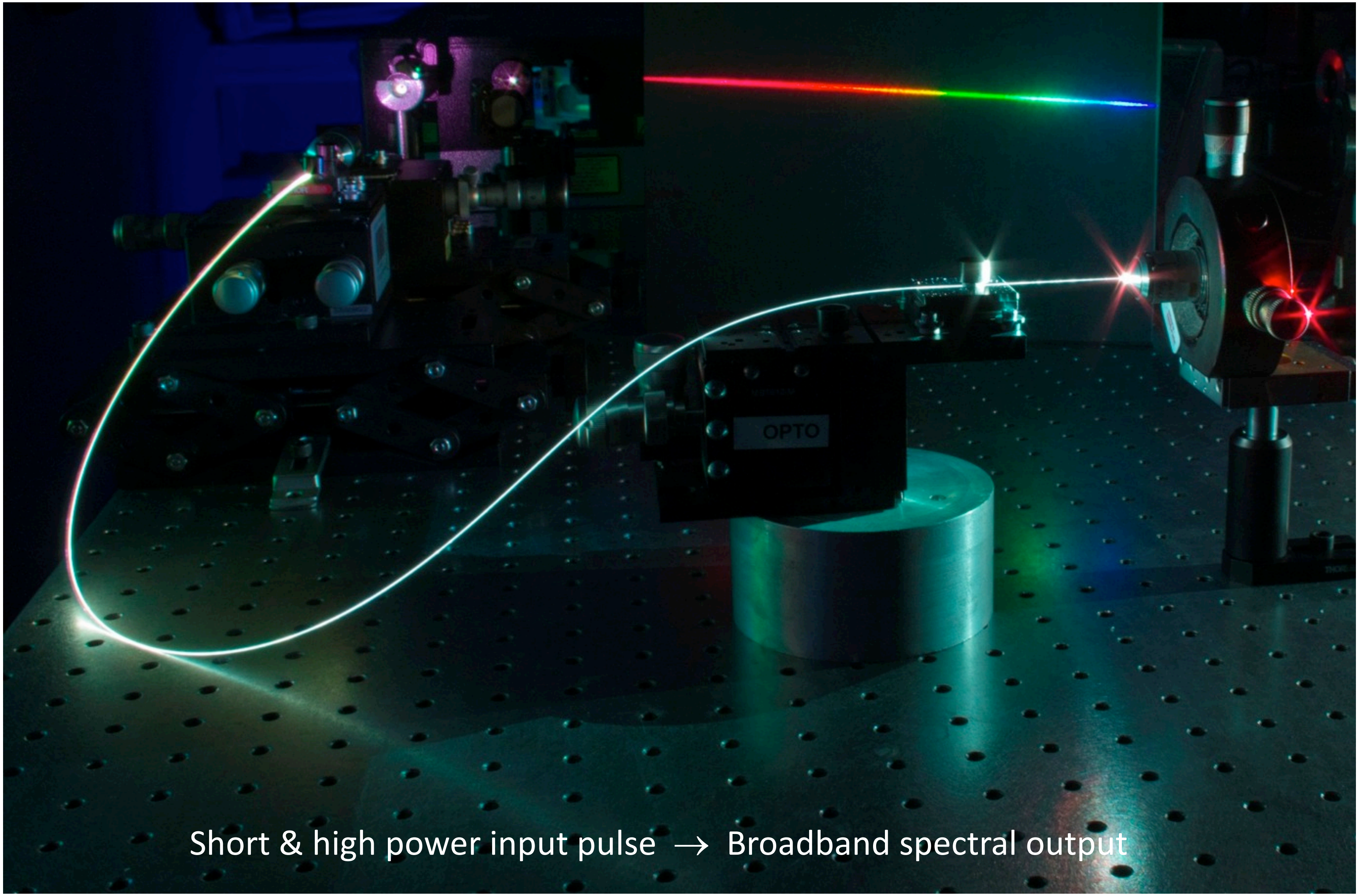
- OPA can provide flat and broad gain bandwidth by dispersion management





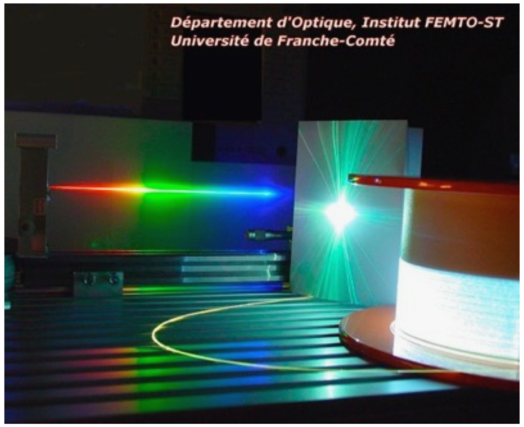


Supercontinuum generation (SC)



Supercontinuum light

- **New hybrid broadband light source: The white-light laser**
- **Broad as the sun (from UV to infrared)**
- **Bright as a laser > 20k times the sun**
- **Spatially coherent - single-mode beam output - Fiber delivery**
- **Temporally coherent to some extent (Soliton Number N<16)**
- **Nothing except synchrotron can give wider bandwidth**



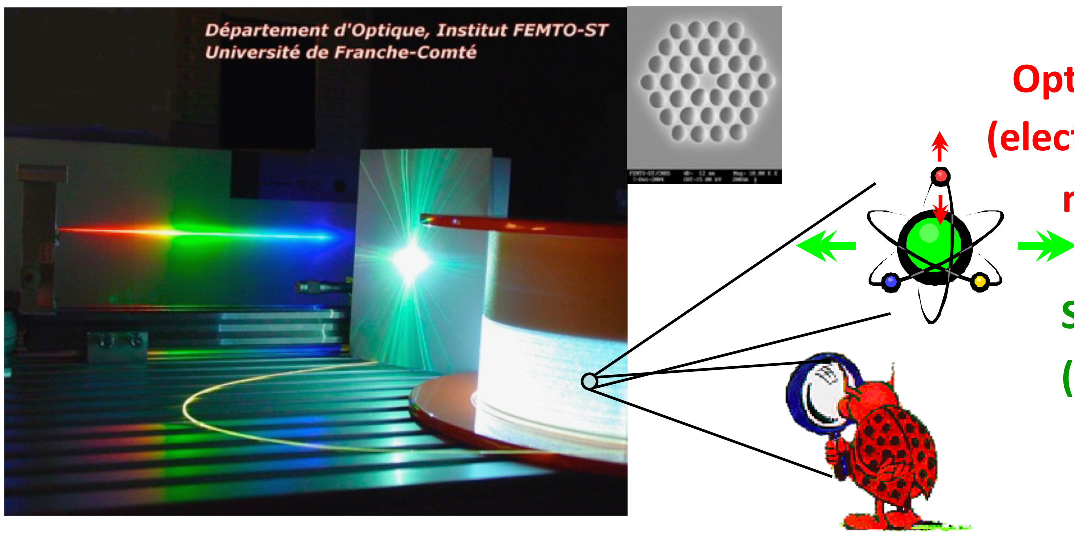
SC generation in a Photonic crystal fiber

NKT Table-top SC source

European synchrotron

Supercontinuum Physics

- **Beautiful fundamental physics !**
- **Exploit fiber nonlinearities to convert laser light to new wavelengths**
- **Based on nonlinear Kerr and Raman responses**
- **Interaction between linear dispersion effects and nonlinear effects**



Optical Kerr Effect
(electronic response)
$n = n_0 + n_2 I(z,t)$

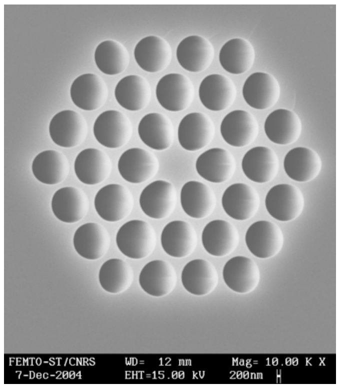
Stimulated Raman Scattering
(Delayed Molecular response)

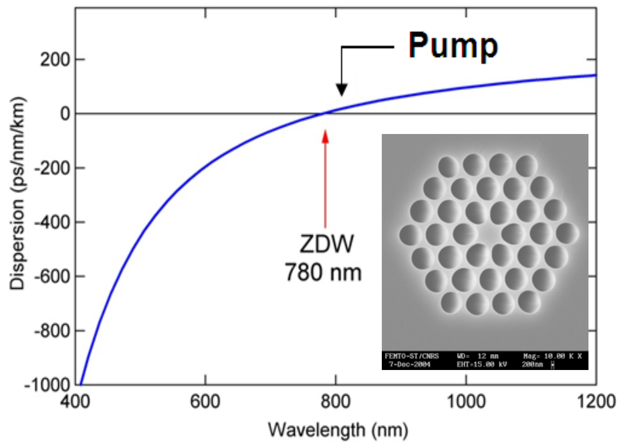
$$P_{NL}(t) = \varepsilon_0 \chi_K^{(3)} : E(t)E(t)E(t) + \varepsilon_0 E(t) \int_{-\infty}^{t} \chi_R^{(3)}(t-t')E(t')E(t')dt'$$

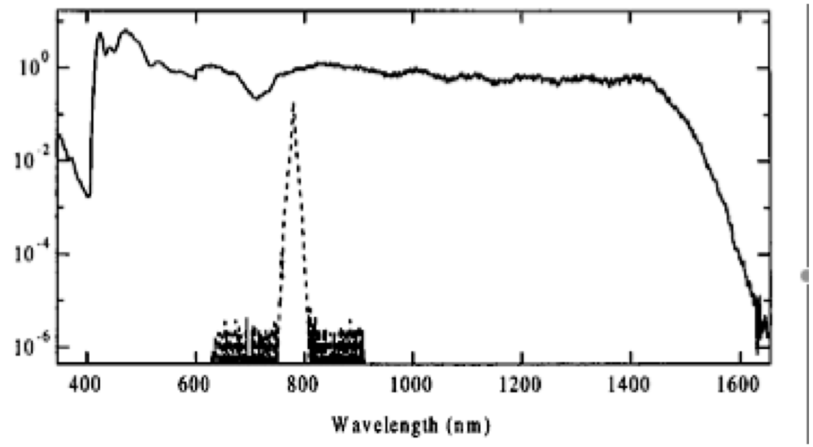
Instantaneous Elastic Kerr effect
(No energy exchange with matter)

Inelastic Raman light scattering:
(Molecular vibrational state)

Supercontinuum Physics







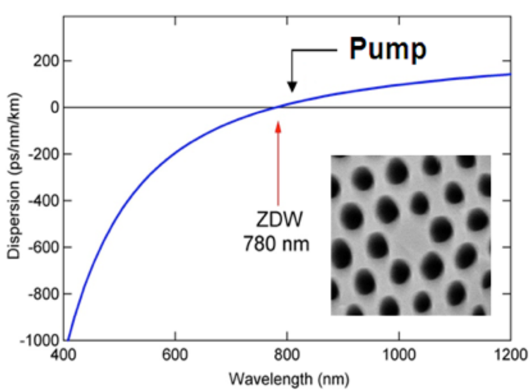
The optical supercontinuum is formed when a high power femtosecond pulse @ 800 nm is sent to a microstructured optical fiber with both small effective mode area and zero-dispersion wavelength close to the pump that cause severe spectral broadening.

Supercontinuum Physics

Modelling with a generalized nonlinear Schrodinger equation (GNLSE)

$$\frac{\partial A}{\partial z} = \underbrace{-\frac{\alpha(\omega)}{2}A}_{\textit{loss}} + \underbrace{\sum_{k\geq 2}\frac{i^{k+1}}{k!}\beta_k\frac{\partial^k A}{\partial T^k}}_{\textit{dispersion}} + i\gamma\underbrace{\left(1 + i\tau_0\frac{\partial}{\partial T}\right)}_{\textit{Self-steepening}}\underbrace{\left(A\int_{-\infty}^{+\infty}R(T')|A(z,T-T')|^2 dT'\right)}_{\textit{Nonlinear response (Kerr + Raman)}}$$

Physics: NLSE + perturbations

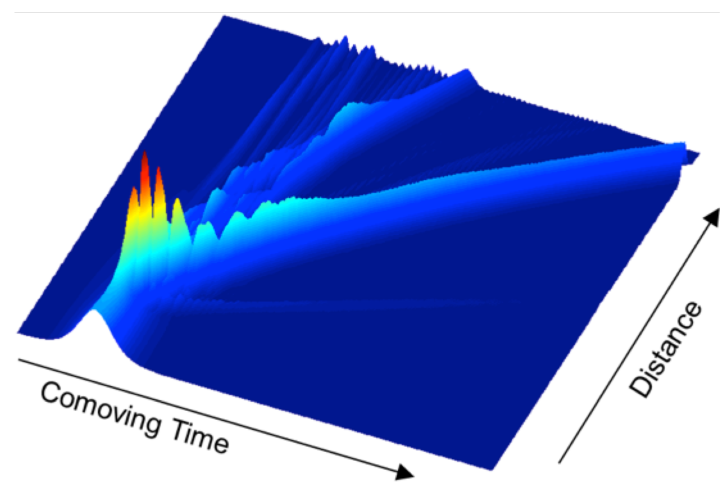


Three main processes

Soliton ejection & fission

Raman – shift to long λ

Radiation – shift to short λ



K. L. Corwin et al., *Phys. Rev. Lett* **90** (2003) J. M. Dudley et al., *Rev. Mod. Phys.* **78** (2006) J. M. Dudley & J. Roy Taylor *Nat. Photon* **3** (2009)

Supercontinuum Physics

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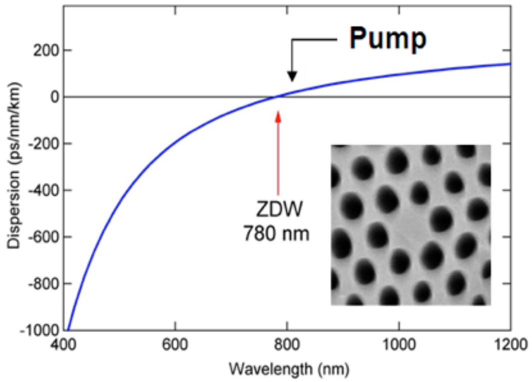
Physics: NLSE + perturbations



Three main processes

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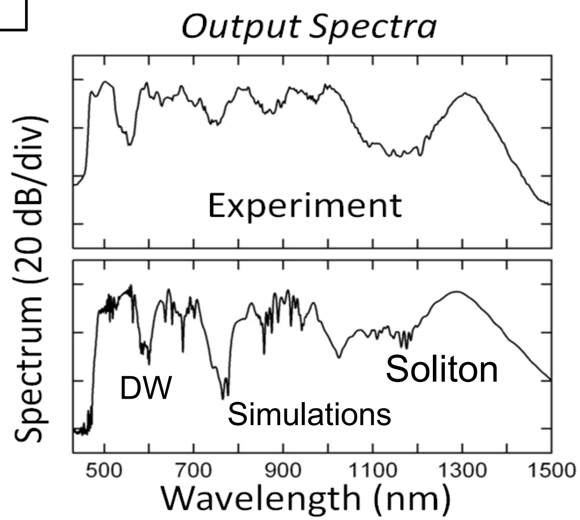
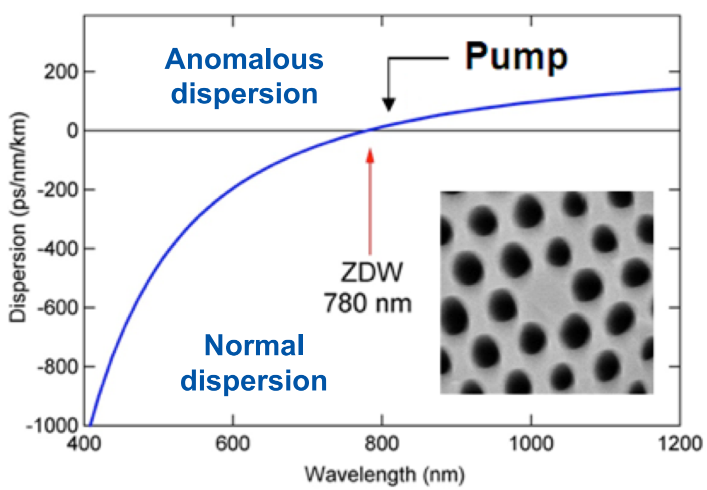


K. L. Corwin et al., *Phys. Rev. Lett* **90** (2003)   J. M. Dudley et al., *Rev. Mod. Phys.* **78** (2006)   J. M. Dudley & J. Roy Taylor *Nat. Photon* **3** (2009)

Supercontinuum Physics : femtosecond regime

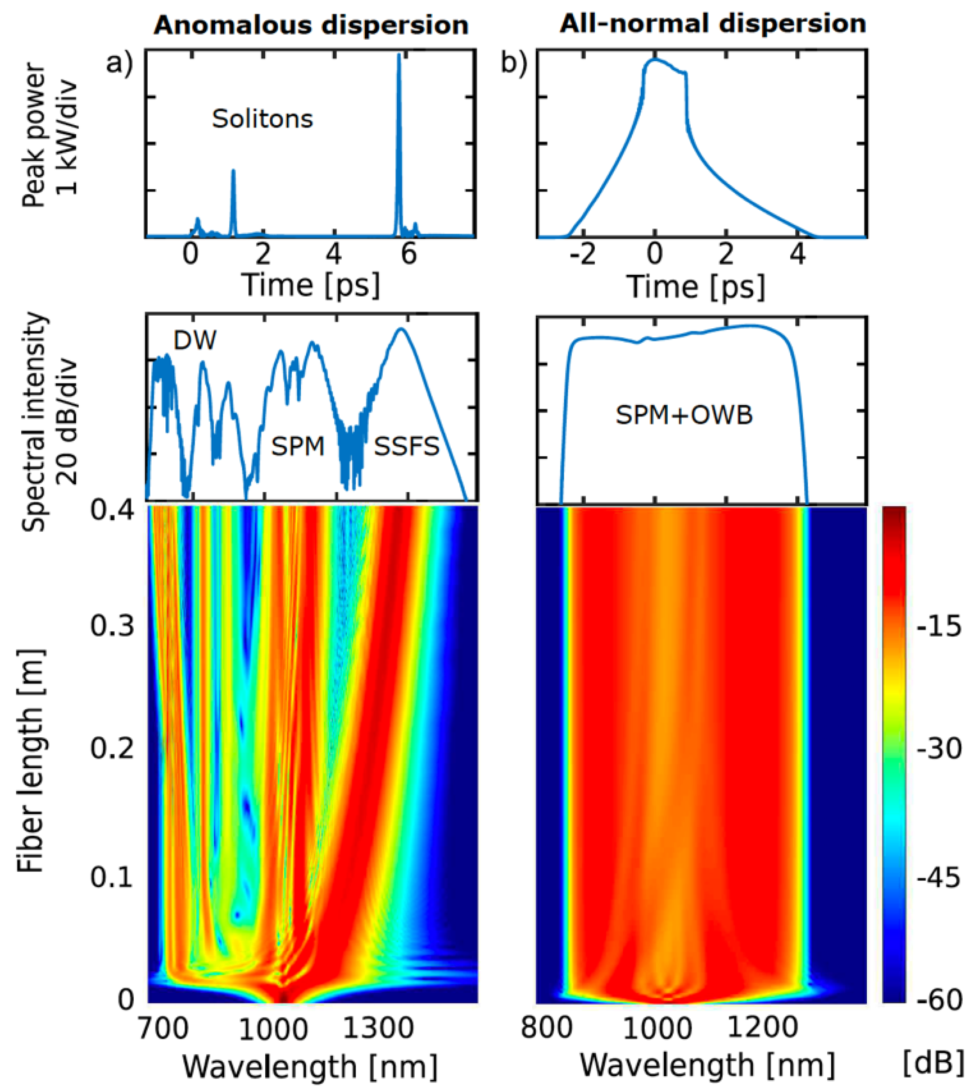
- **Anomalous dispersion (D>0, β_2<0):**
- Soliton Formation & Fission,
- Raman shift (SSFS) to long l
- Dispersive-wave (DW) to short l

- **Normal dispersion (D<0, β_2>0):**
- Self-phase modulation (SPM)
- Optical-wave breaking (OWB)



J. M. Dudley et al., Rev. Mod. Phys. ***78*** *(2006)*

$$\frac{\partial A}{\partial z}+\frac{\alpha}{2}A-\sum_{k\geq 2}\frac{i^{k+1}}{k!}\beta^k\frac{\partial^k A}{\partial T^k}=i\gamma\left(1+i\tau_{shock}\frac{\partial}{\partial T}\right)\left(A(z,T)\int_{-\infty}^{+\infty}R(T')\,|A(z,T-T')|^2\,dT'\right)$$



Dispersive Wave (resonant radiation)

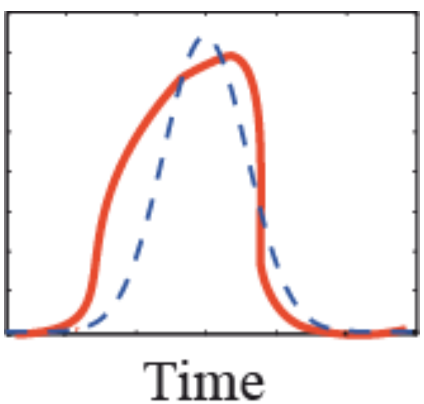
Impact of third-order dispersion (β_3) on solitons

PHYSICAL REVIEW A VOLUME 51, NUMBER 3 MARCH 1995

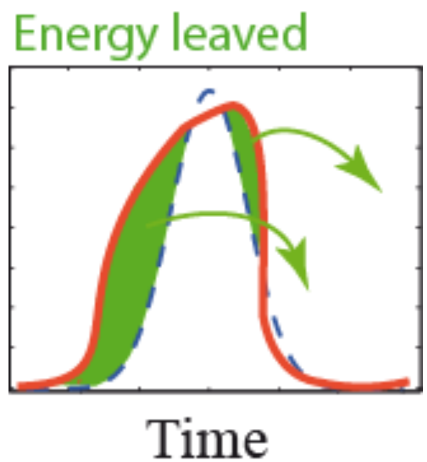
Cherenkov radiation emitted by solitons in optical fibers

Nail Akhmediev and Magnus Karlsson*

Soliton instability due to the Third-order dispersion or Raman scattering



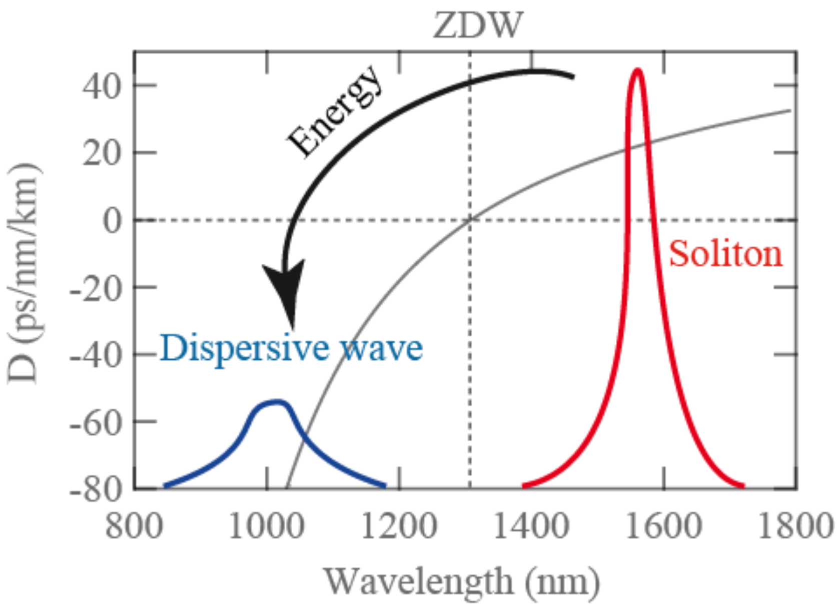
Energy dissipation to become stable



But where energy is leaking ?

Frequency-shifted Dispersive waves

Dispersive Wave (resonant radiation)

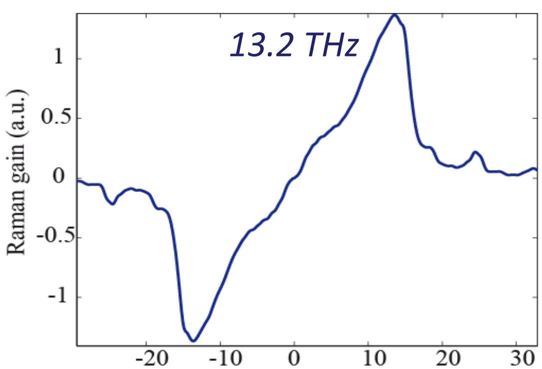


Phase-matching condition between solitons and dispersive waves yields:

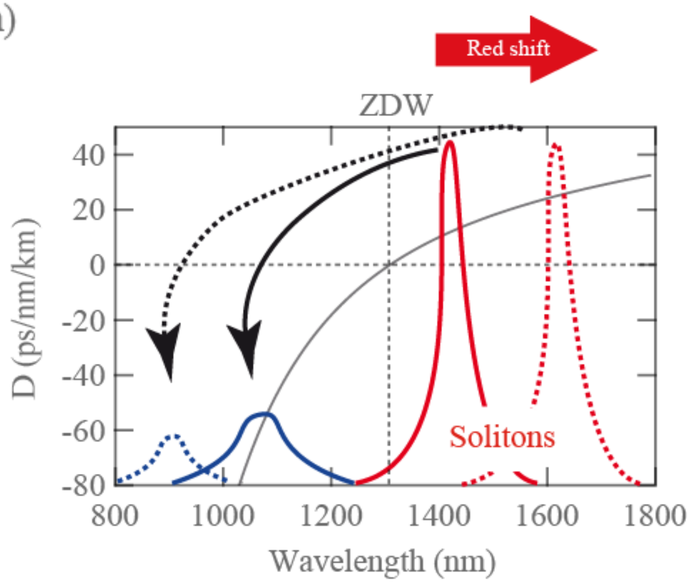
$$\beta(\omega_{DW}) - \beta(\omega_S) - \beta(\omega_{DW} - \omega_S) - \gamma P_s/2 = 0$$

$$\Delta\omega_{Disp.Wave} \approx -\frac{3\beta_2(\omega_{sol.})}{\beta_3(\omega_{sol.})} + \frac{\gamma P_{Maxsol}\beta_3(\omega_{sol.})}{3\beta_2^2(\omega_{sol.})}$$

Raman gain spectrum



wavelength



SC red side: soliton self-frequency shift (SSFS)

SC blue side: Dispersive waves trapped by solitons

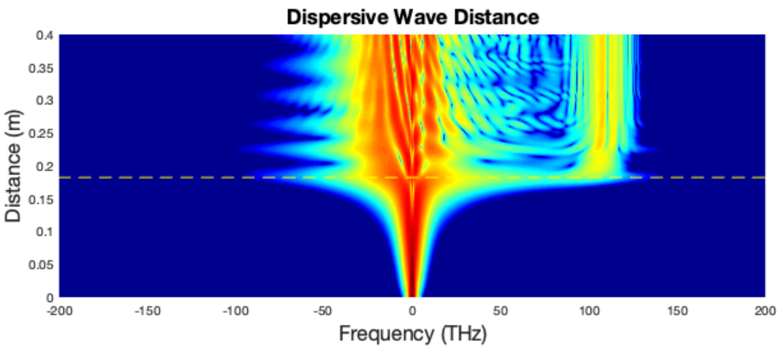
Dispersive Wave (resonant radiation)

Dispersive wave distance emitted from solitons due to β_3 or β_4

β_2 = -8e-27 s^2/m; β_3=4e-41 s^3/m; β_4=0

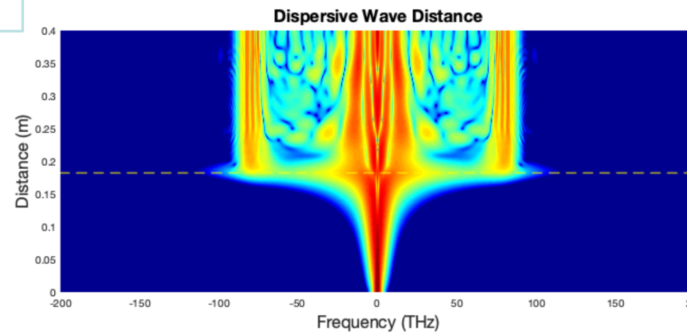
P_0=1 kW
T_0=200 fs
γ=27 $W^{-1}km^{-1}$

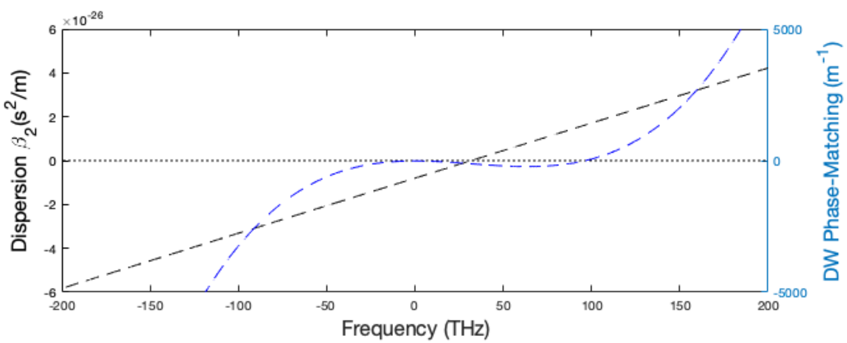
β_2 = -8e-27 s^2/m; β_3= 0; β_4=4e-41 s^4/m

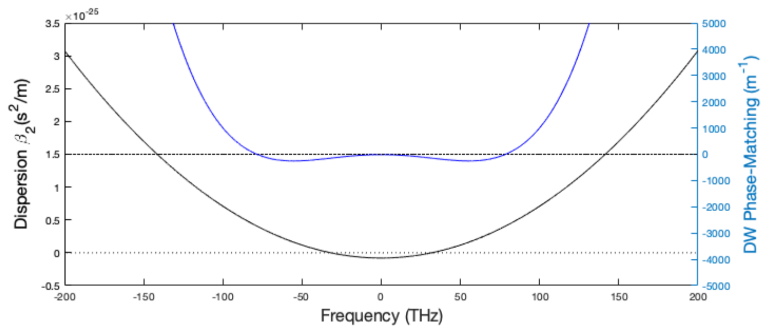


$$L_{WB} = \sqrt{\frac{3/2}{N^2+1}}$$

$$N = \sqrt{\frac{\gamma P_0 T_0^2}{|\beta_2|}}$$



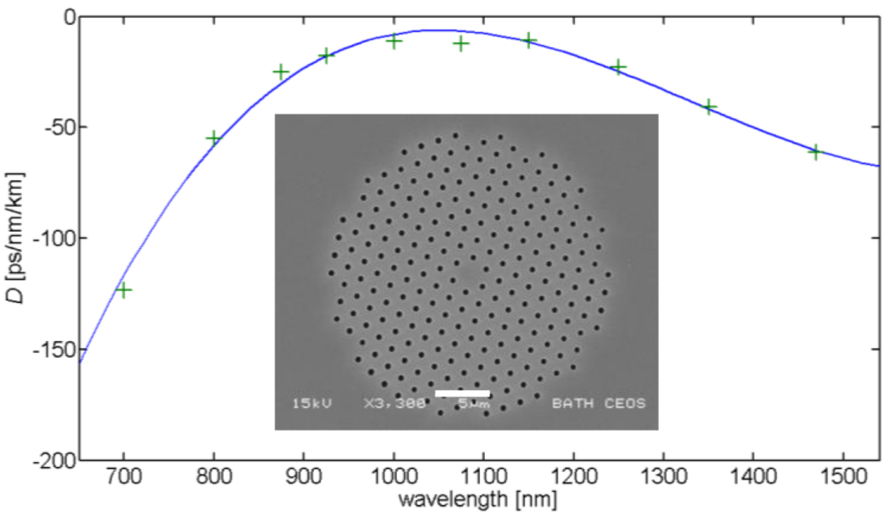




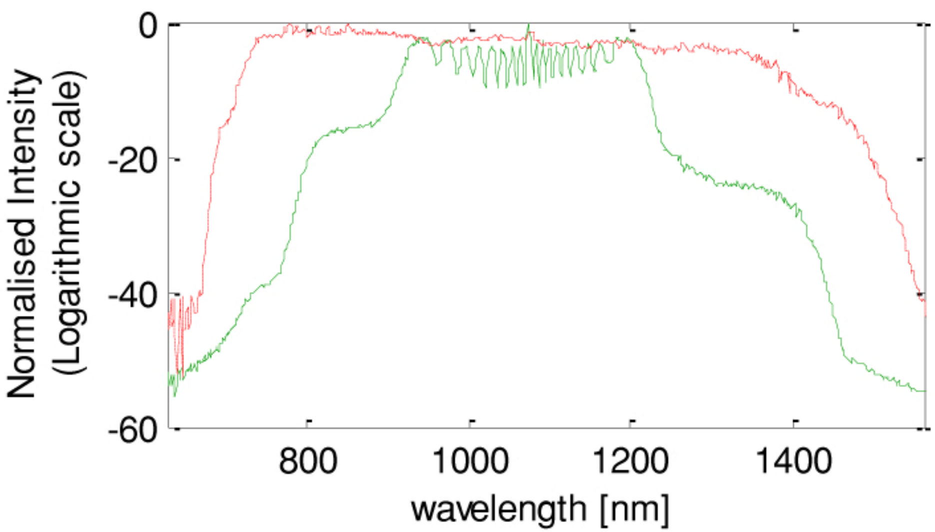
Coherent Supercontinuum generation

- Pumping in the low normal dispersion regime
- Self-phase modulation (SPM) + Optical-Wave Breaking (OWB)
- OWB = Dispersive Shock Wave (DSW)
- Smooth and flat SC spectrum – Pulse-preserved
- Gives low noise and high coherence level

PCF with all-normal group velocity dispersion



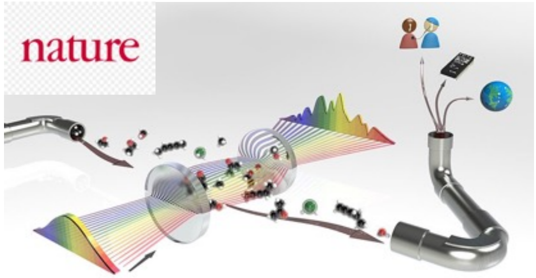
Wideband supercontinuum generation



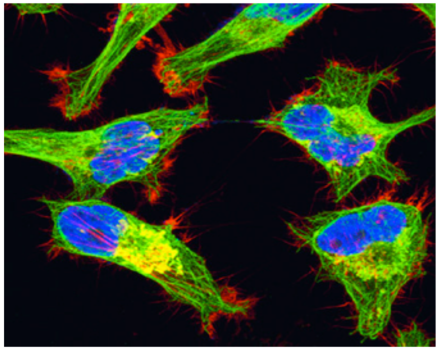
L. E. Hooper, P. J. Mosley, A. C. Muir, W. J. Wadsworth, J. C. Knight, Opt. Express **19,** 4902 (2011).

Supercontinuum Applications

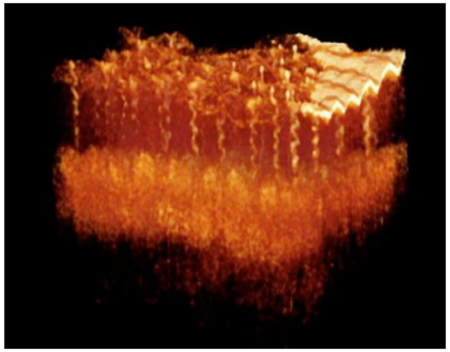
- Biomedical imaging and microscopy,
- OCT, Confocal, Fluorescence, CARS imaging
- Cell-Tissue-Material analysis, Air pollution
- Molecular spectroscopy, Material processing
- Remote non-destructive detection
- Optical Frequency Comb (OFC) Metrology
- LIDAR, flow cytometry, and many others…



Molecular spectroscopy



Fluorescence imaging



OCT image of a finger tip

TheScientist TOP INNOVATIONS OF 2008

Confocal microscopy

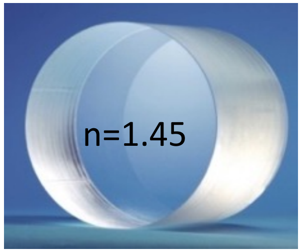
SUPUVIR: SC Fiber materials

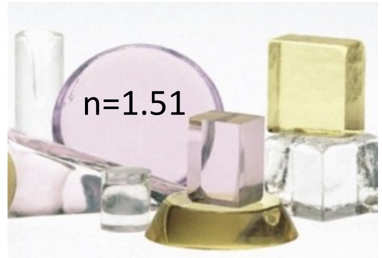
Silica glasses	Fluoride glasses	Chalcogenide glasses

UV-grade glass

High-OH dopant

Gas-filled fiber

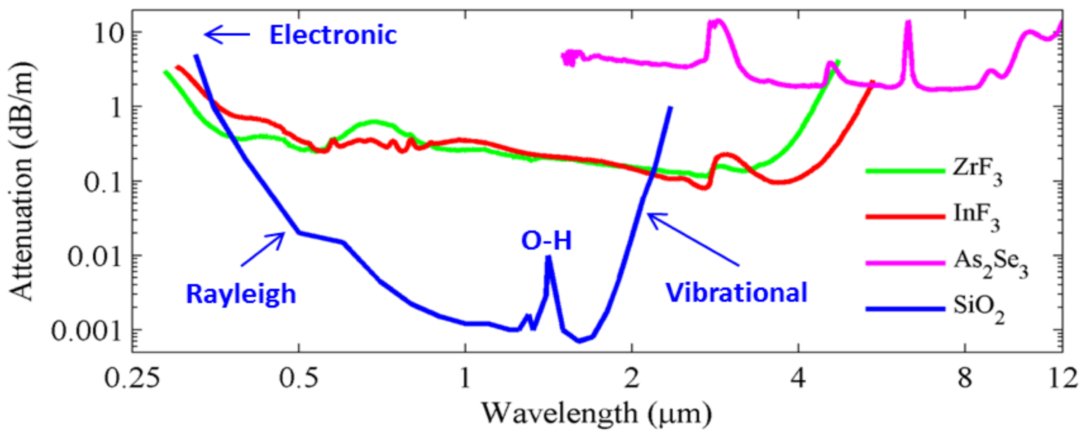


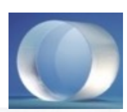


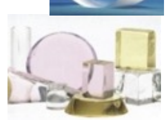
Heavy-metal-oxyde (HMO) glasses

Tellurite glasses (TZN)

Telluride glasses (Te)

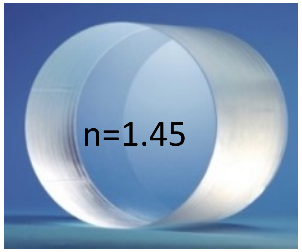


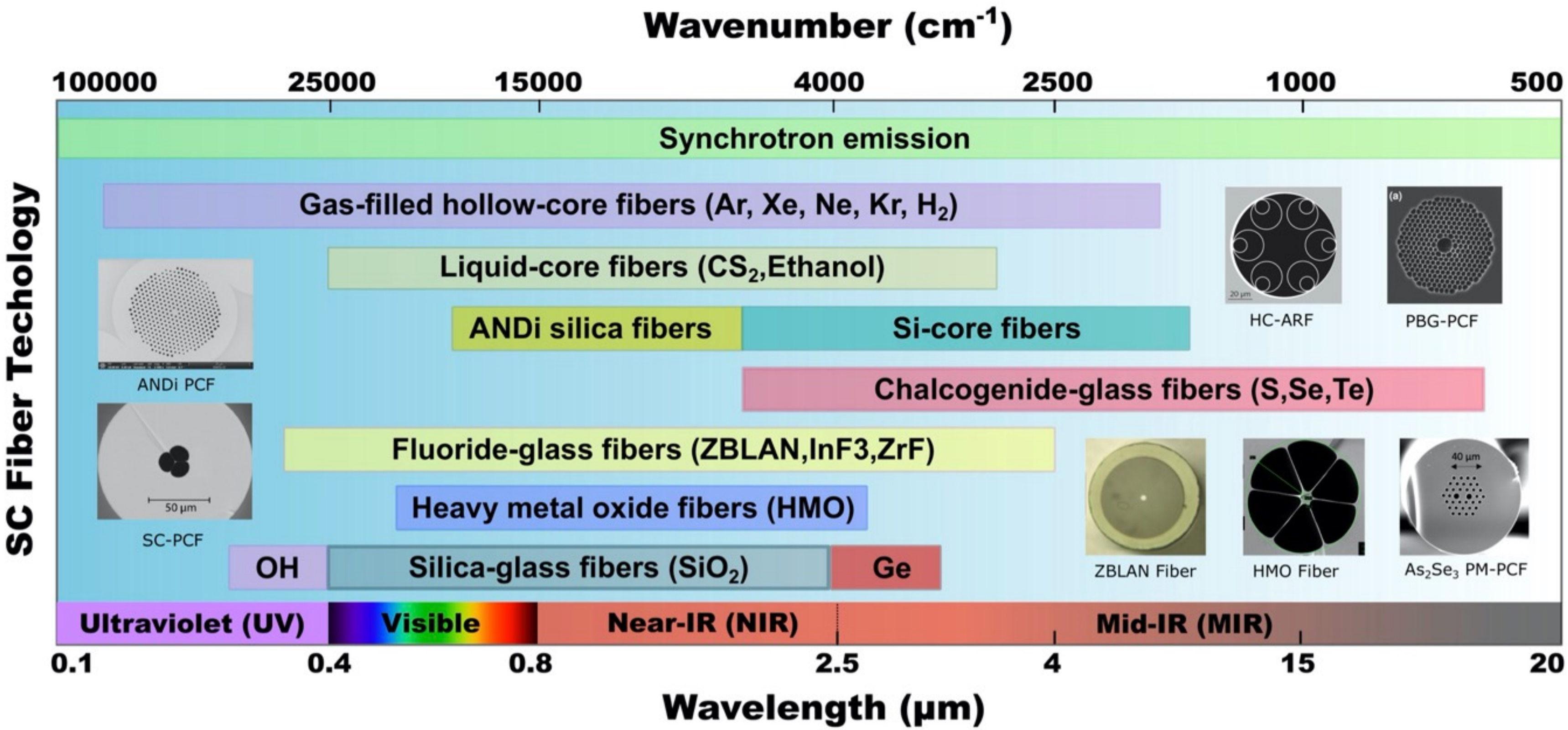




Glass	n	$n_2\,(10^{-20}\mathrm{m^2W^{-1}})$	T_g (°C)	ZDW (μm)
Silica	1.45	2.7	~1200	1.26
ZBLAN	1.48-1.53	2.1-2.55	230-300	1.62-1.71
$As_{40}Se_{60}$	2.82-2.9	1400-3000	178 -185	7.2-7.4

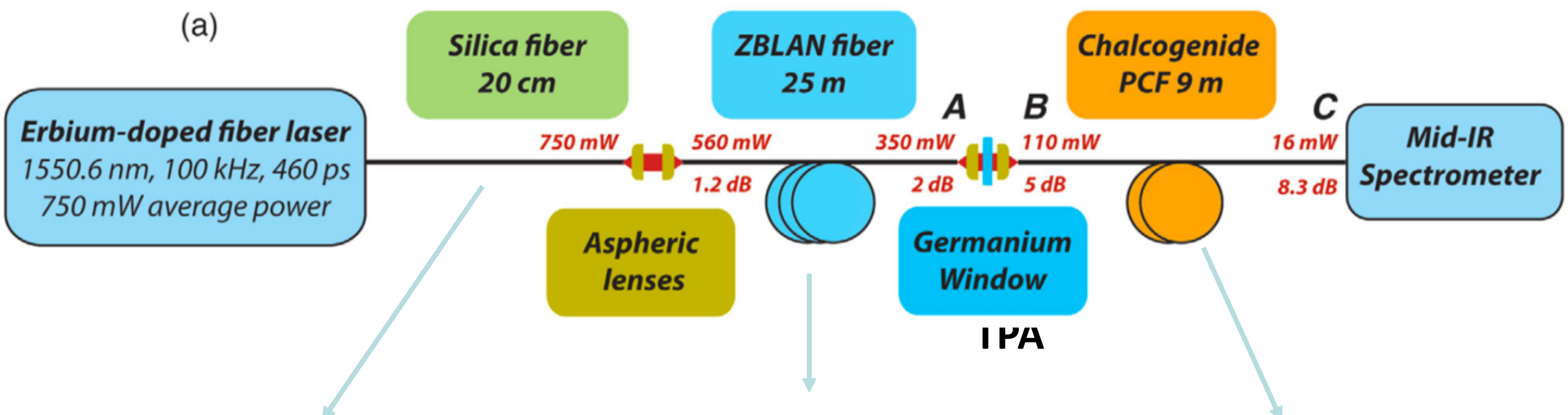
SUPUVIR: SC Fiber materials

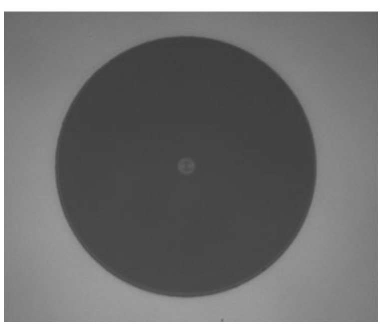
	Silica glasses	Fluoride glasses	Chalcogenide glasses	
UV-grade glass High-OH dopant Gas-filled fiber				Heavy-metal-oxyde (HMO) glasses Tellurite glasses (TZN) Telluride glasses (Te)



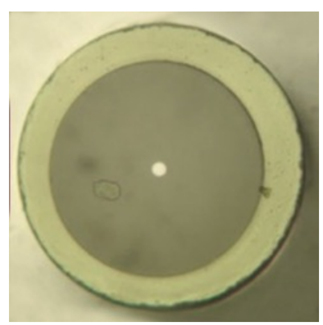
Cascaded Mid-IR SC Source

- **Development of compact cascaded fiber mid-IR SC light sources**
- Combine silica - fluoride - chalcogenide fibers
- Directly pumped by a table-top fiber laser at 1.55 µm

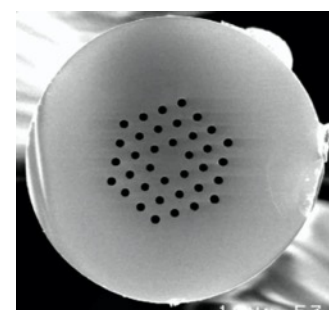




Silica SMF-28



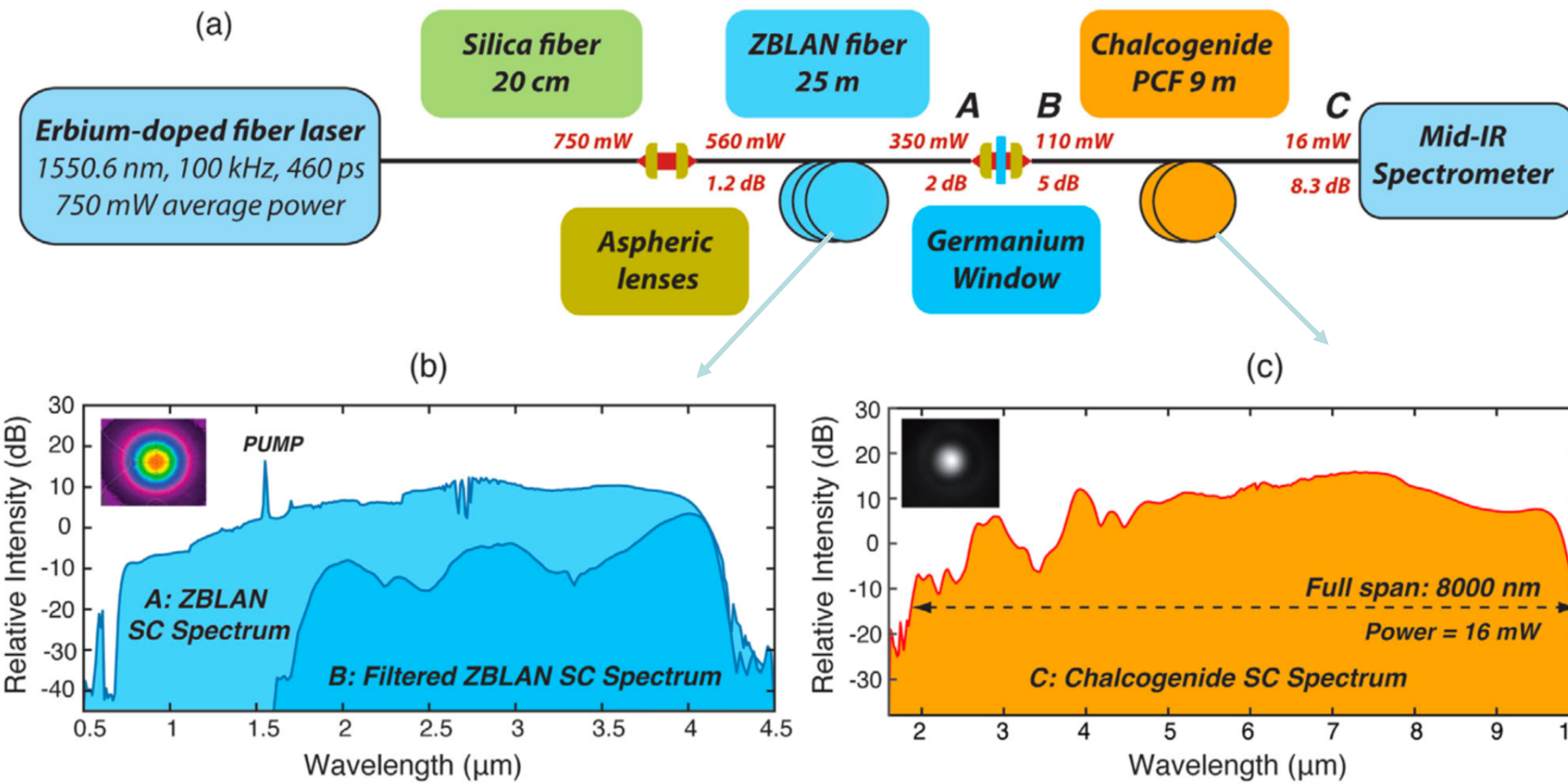
ZBLAN step-index single-mode fiber



$As_{38}Se_{62}$ PCF

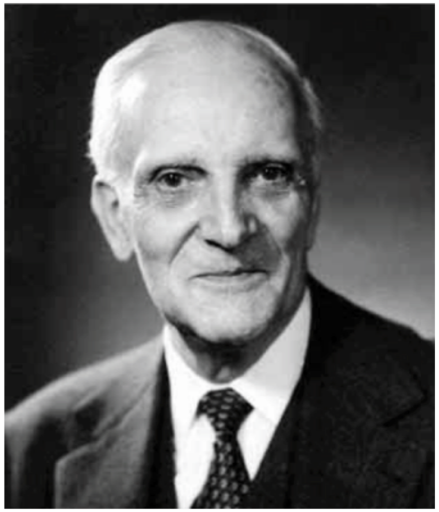
Cascaded Mid-IR SC Source

- **Development of compact cascaded fiber mid-IR SC light sources**
- Step-wise SC extension towards the MIR by Raman soliton drifting
- 2-10 µm wide SC spectrum with 16 mW output power



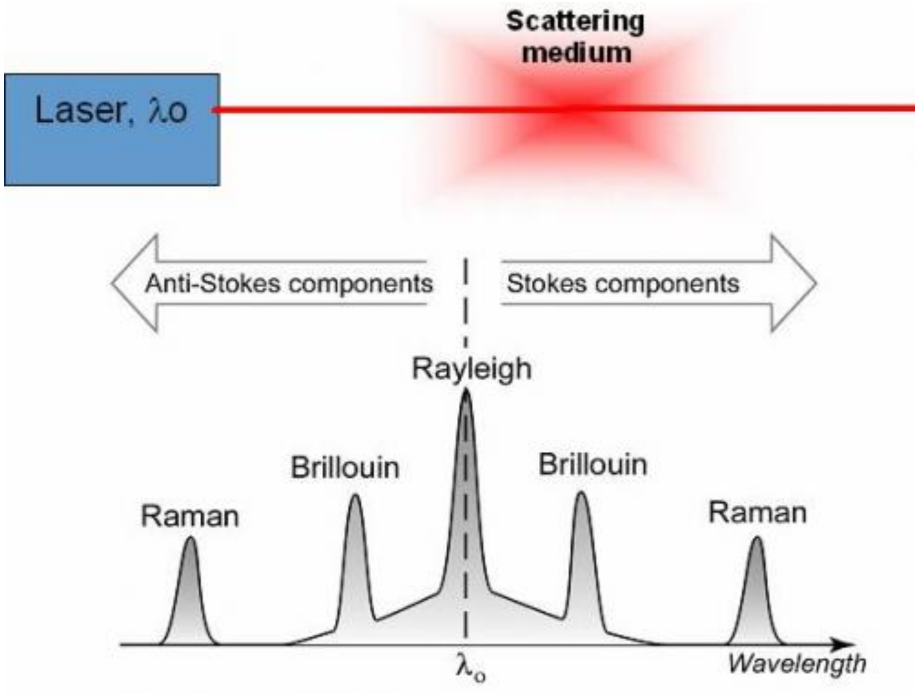
S. Venck et al, Laser and Photonics Rev. 14 (6), 2000011 (2020)

Brillouin scattering in optical fibers



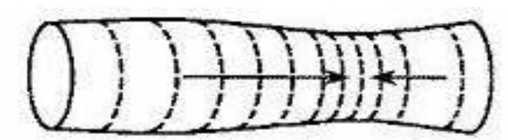
Léon Brillouin (1889-1969)

- Inelastic scattering by hypersound acoustic waves
- Brillouin scattering is a photon-phonon interaction



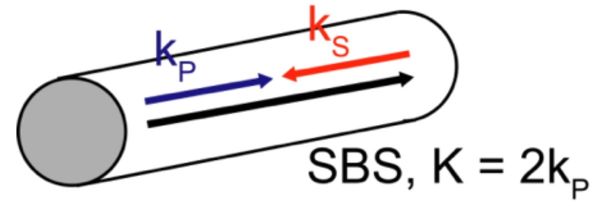
Brillouin scattering in optical fibers

Longitudinal wave



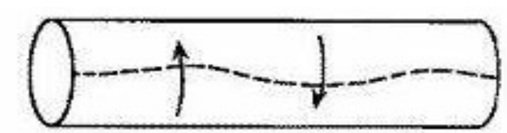
Stimulated Brillouin Scattering (SBS)
Velocity=5960 m/s
Wavelength ~ 0.5 µm
Frequency = 11 GHz

E. Ippen. & R. Stolen *Appl. Phys. Lett.* **21** (1972)



Backward scattering

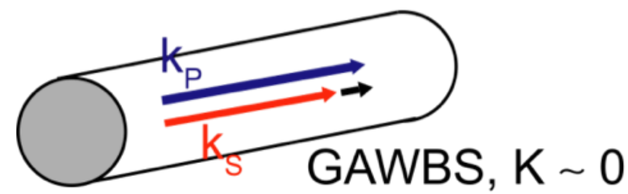
Torsional wave



Radial wave

Guided acoustic wave Brillouin Scattering (GAWBS)
Shear velocity=3600 m/s
Wavelength ~ 0.5 µm
Frequency [100 MHz-1 GHz]

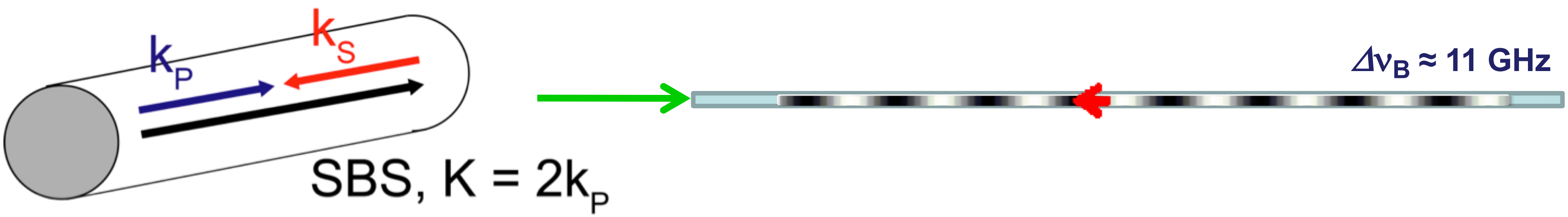
R. M. Shelby et al. *Phys. Rev. B* **31**, 5244 (1985)



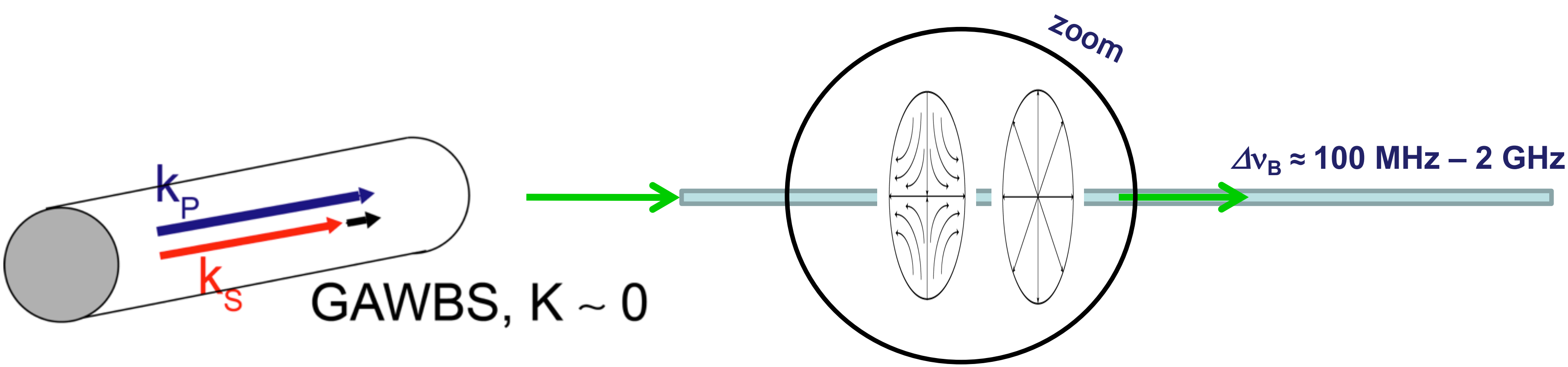
Forward scattering

Brillouin scattering in optical fibers

Backward Stimulated Brillouin scattering (SBS): elastic longitudinal wave (acoustic phonon)

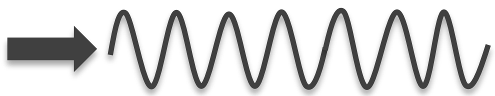


Forward Brillouin scattering: Transverse acoustic wave (Guided Acoustic Wave Brillouin Scattering)

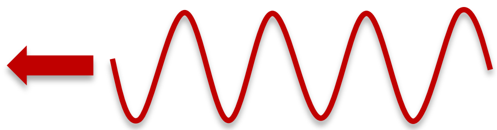


Brillouin scattering in optical fibers

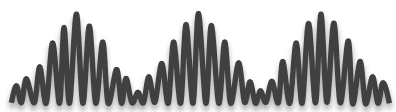
(1) Incident pump beam



(2) Spontaneous scattering of light by acoustic phonons

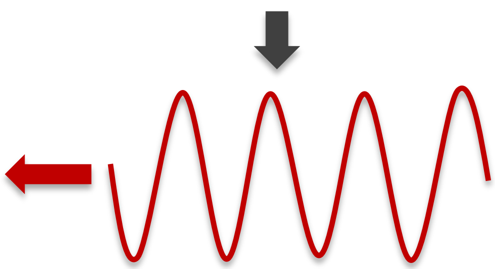


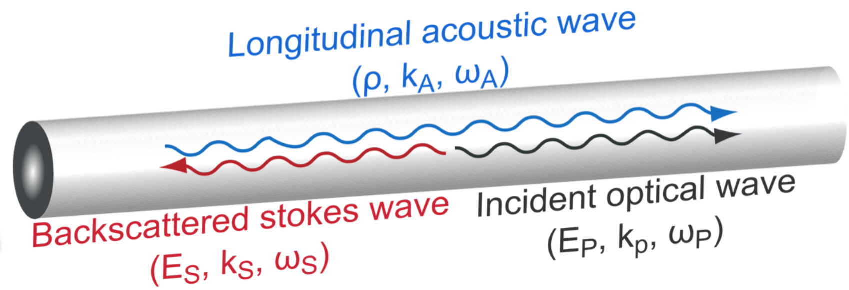
$|(1)+(2)|^2 \rightarrow$ (3) Interference



(3) → (4) Acoustic wave created by electrostriction

Exaltation of light scattering: stimulated scattering





Phase matching condition

$$k_B = k_P - k_S$$

Conservation of energy

$$\omega_B = \omega_P - \omega_S$$

Brillouin gain

$$g_B(\omega) = \frac{g_{B0}(\Delta\omega_B/2)^2}{(\omega-\omega_B)^2 + (\Delta\omega_B/2)^2}$$

Brillouin frequency shift (Doppler effect)

$$\nu_B = \frac{2 v_A n_{eff}}{\lambda}$$

Silica $v_A = 5960\, m.s^{-1}$ $\nu_B \approx 11 GHz$ $n_{eff} \approx 1{,}44$

R. W. Boyd, *Nonlinear Optics,* third edition

Brillouin scattering in optical fibers

Pump and Probe at slightly different frequencies ν_P and ν_S are interacting through acoustic wave at Brillouin frequency ν_B. ($\nu_P=\nu_S+\nu_B$ energy conservation)

Moving index Bragg grating generated along the fiber core by electrostriction

CW pump (strong)

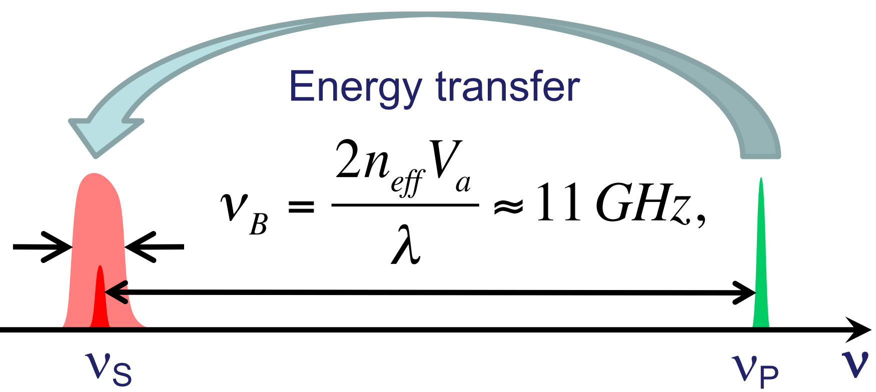
CW probe @ Brillouin shift

$$V_a \approx 6000\, m.s^{-1}$$

SBS gain process on the Stokes side

$\Delta\nu_B$ ~ 26 MHz in silica fiber

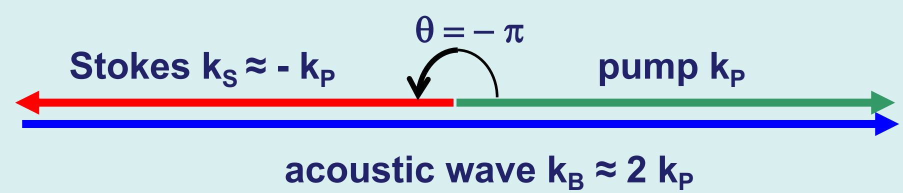
acoustic lifetime τ_a ~ 10 ns



$$\nu_B = \frac{2 n_{eff} V_a}{\lambda} \approx 11\, GHz,$$

Phase-matching condition:

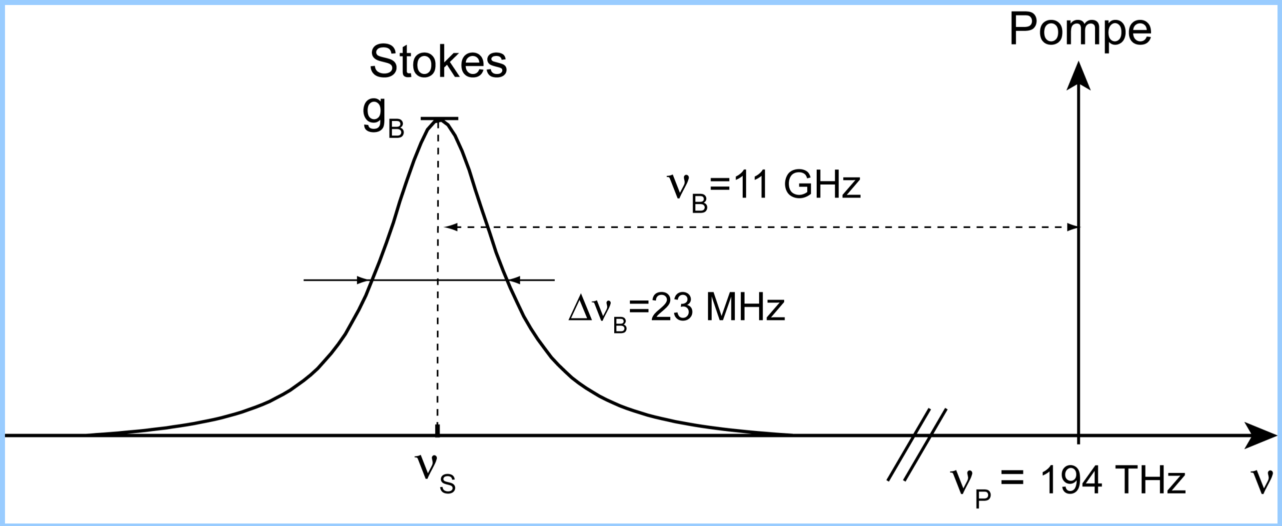
$$\overrightarrow{k_P} = \overrightarrow{k_S} + \overrightarrow{k_B}$$



* Starts from thermal agitation, i.e. spontaneous scattering, when no probe

Brillouin scattering in optical fibers

In silica fiber, @ 1550 nm, Brillouin gain = $3.2.10^{-11}$ m/W



➢ Brillouin frequency

$$\nu_B = \frac{2\, n\, v_a}{\lambda_P}$$

➢ Brillouin spectrum

$$g_B(\nu) = \frac{(\Delta\nu_B/2)^2}{(\nu - \nu_B)^2 + (\Delta\nu_B/2)^2}\, g_B$$

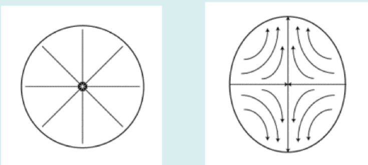
➢ Brillouin gain at the resonance

$$g_B = \frac{2\,\pi\, n^7\, p_{12}^2}{c\, \lambda_P^2\, \rho_0\, v_A\, \Delta\nu_B}$$

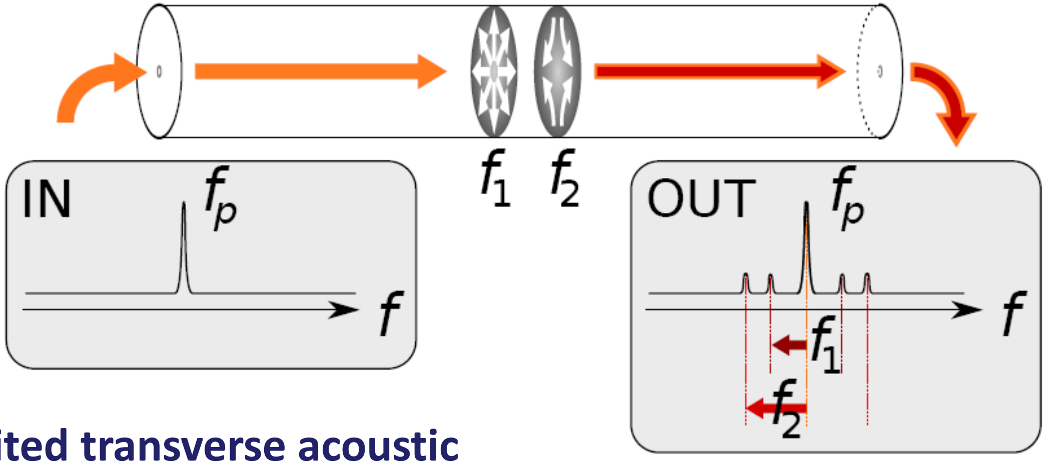
Forward Brillouin scattering in optical fibers

Forward scattering by transversely trapped acoustic resonances

radial & torso-radial modes



polarized & depolarized GAWBS

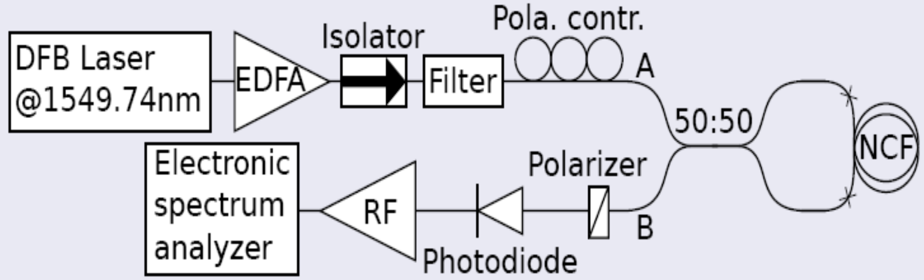


$f_{1,2} \sim$ 100 MHz - 2 GHz
(backscattering : $\sim$ 11 GHz)

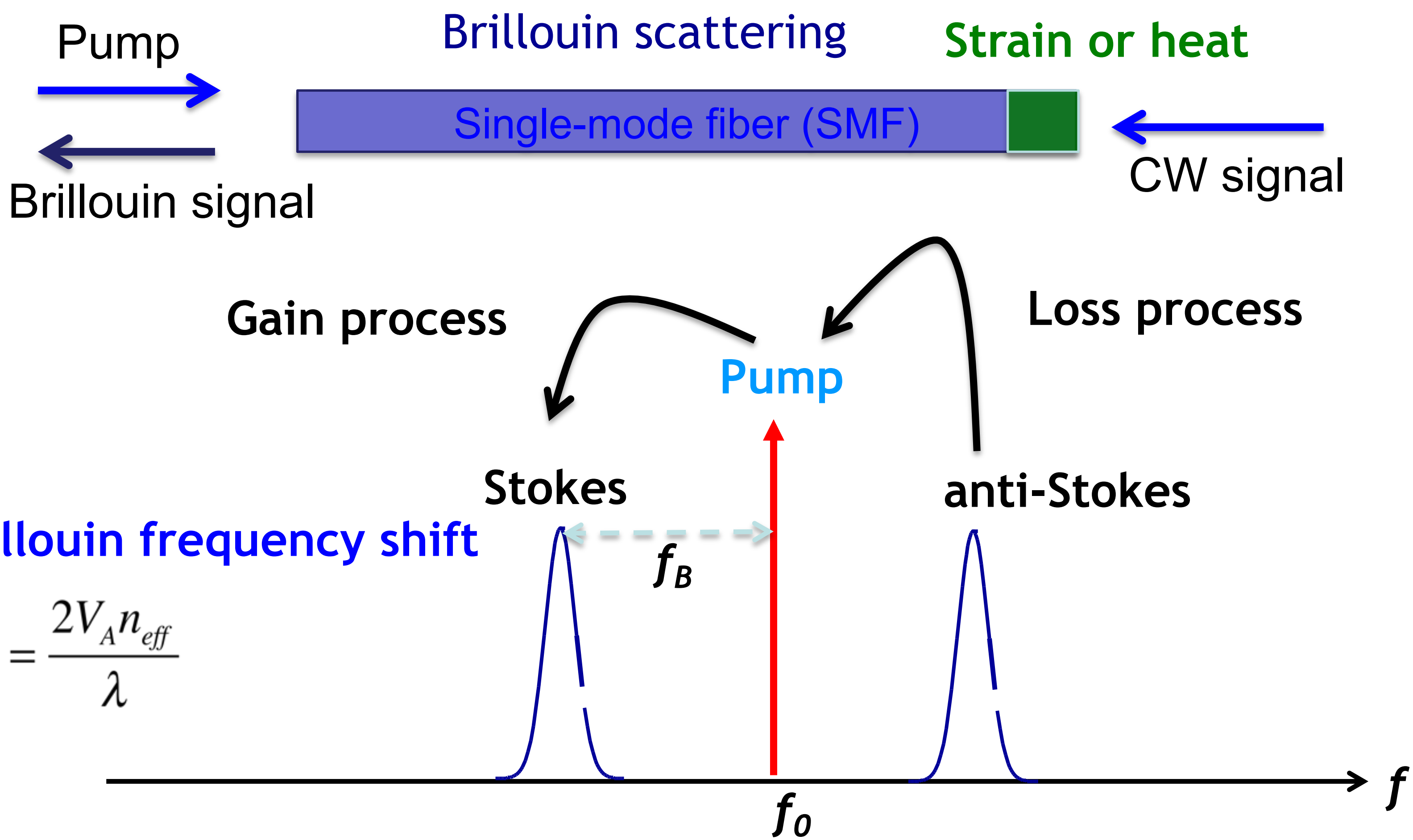
- **Originates from thermally excited transverse acoustic phonon of the fiber (the whole cylinder)**

- **Acoustic modes alter the refractive index periodically**
  ⇨ **phase or polarization modulation of the light wave**

Setup



Stimulated Brillouin scattering (SBS)



f_B **: Brillouin frequency shift**

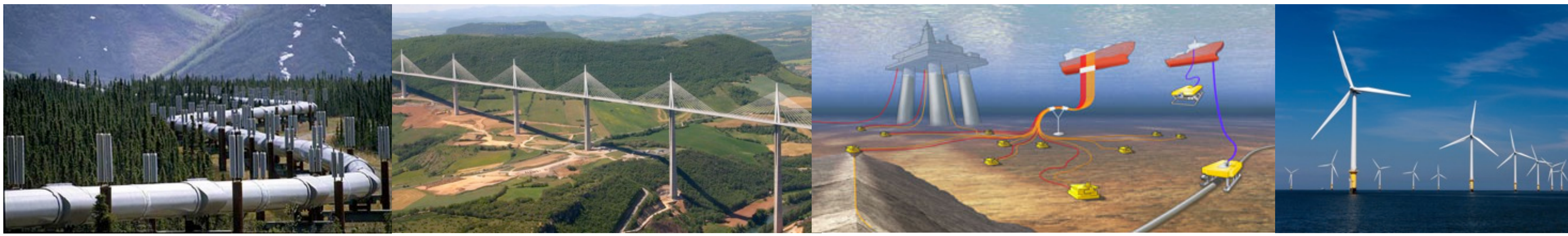
$$f_B = \frac{2V_A n_{eff}}{\lambda}$$

$f_B(\varepsilon,T)= f_B + C_1*\varepsilon + C_2*T$

C_1=0.05 MHz/με, C_2=1 MHz/°C

Brillouin-based applications

Applications to civil engineering and petroleum industry



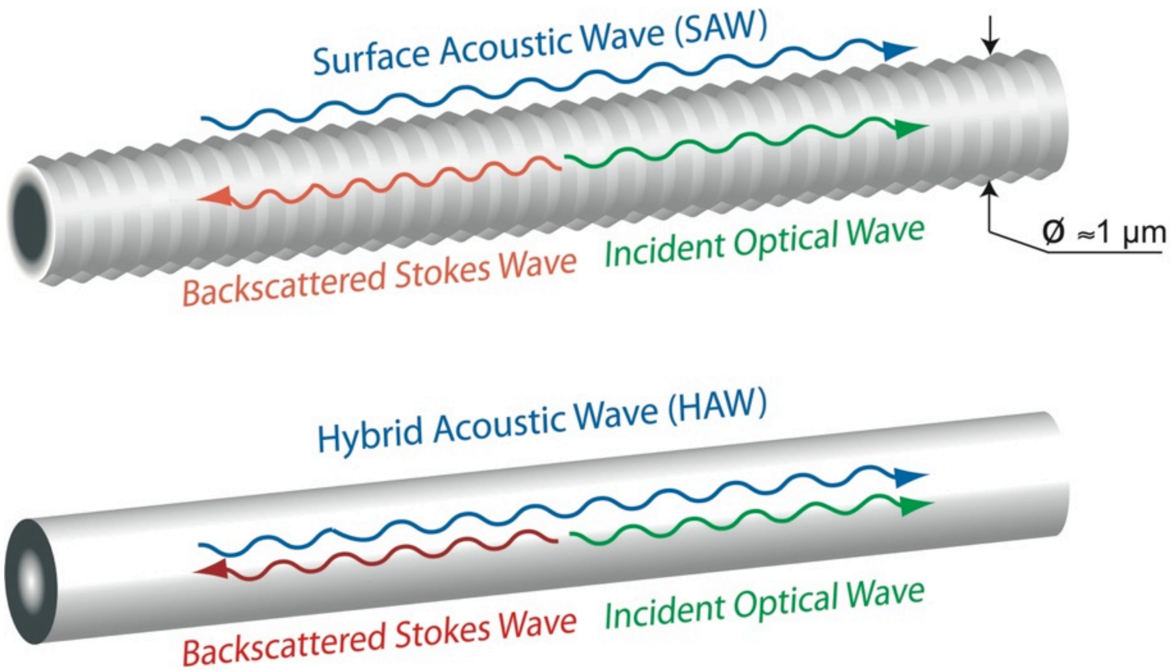
Pipeline (~ 100 km) | Millau Bridge (2460 m) | Petroleum platform | Aeolian

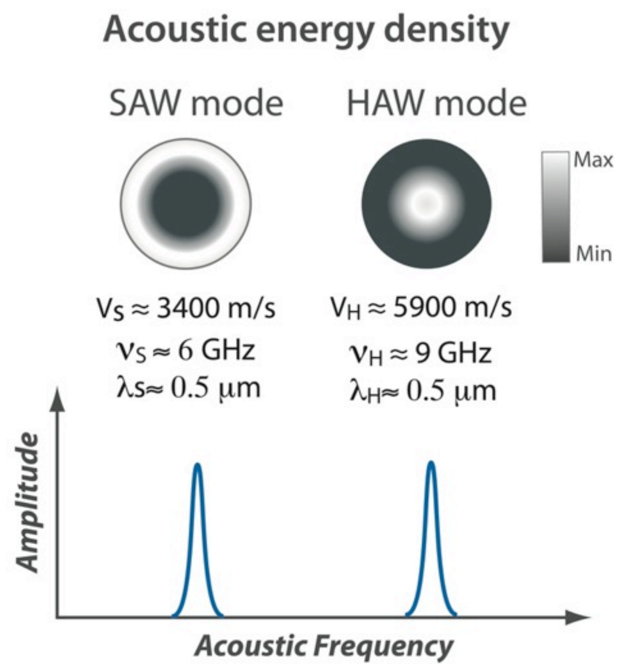
omnisens
Securing asset integrity

Brillouin optical time domain analysis sensors
(Current Performances : Range =30 km; Resolution=1m)

Brillouin scattering in a fiber taper

Unlike optical fibres, optical microwires can actually exhibit both surface acoustic waves (SAW) and hybrid acoustic waves (HAW).





Surface waves travel at a velocity of 3400 m/s leading to new Brillouin lines frequency-shifted by 6 GHz in the backscattered spectrum

Conclusions and Outlooks

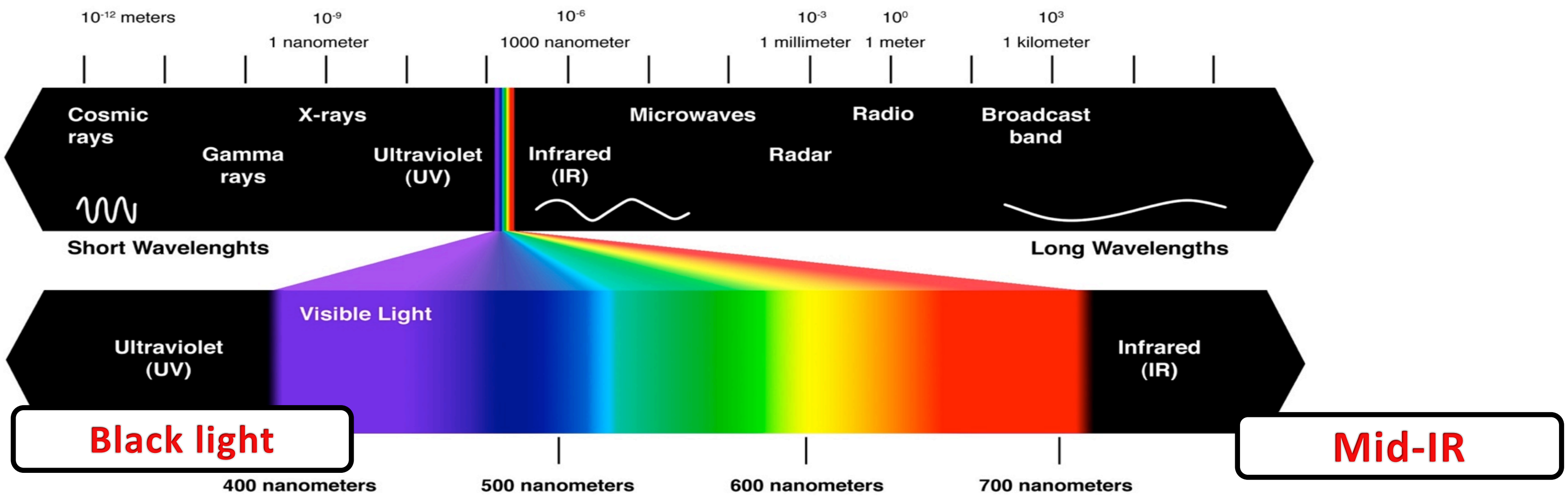
- ✓ Nonlinear optical effects are remarkable physical phenomena that offer unique and exciting potential applications, such as devices in which light can be controlled by light.

- ✓ Due to their strong light confinement capabilities, optical fibers have been early recognized as an ideal medium to exploit the nonlinear effects for all-optical processing applications in telecommunication and fiber laser industries, supercontinuum sources, etc…

- ✓ This research field has recently been stimulated by the development of a new generation of highly nonlinear tiny optical waveguides, e.g., photonic crystal fibers, micro and nanowires, photonic chip, microresonators.

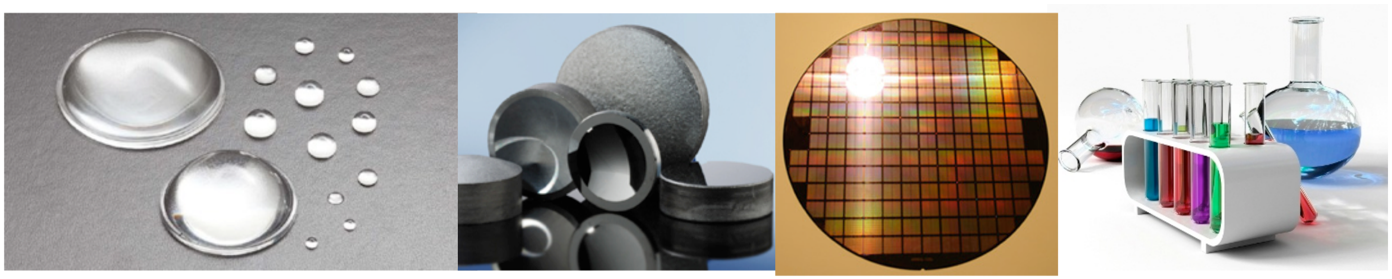
What next ?

Extending applications to UV & Mid-IR wavelength ranges



Exploiting new nonlinear materials

UV Silica, Chalcogenide, Tellurite, Silicium, Liquids, Gas, Crystal fibres ($LibNO_3$)



What next ?

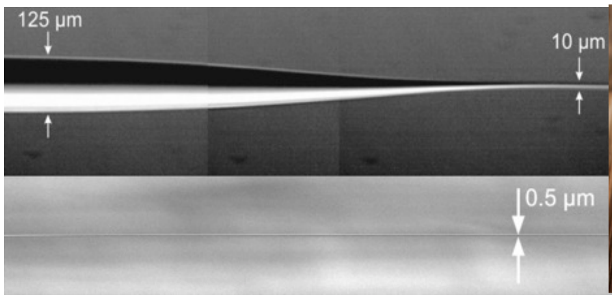
Reducing the dimensions (sub-λ) and the power levels

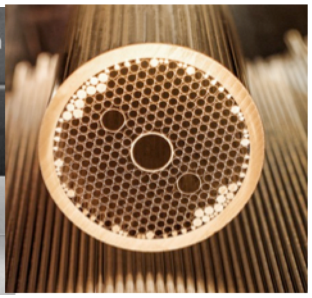
Nanowires, Tapers

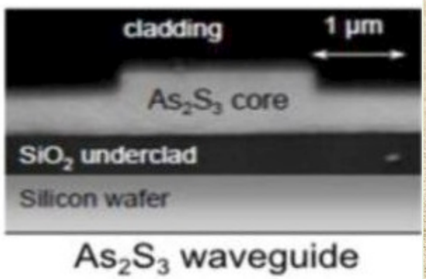
Nanostructured fibers

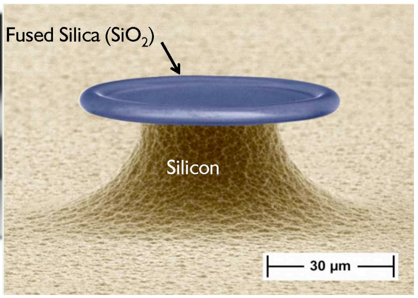
Chip waveguides

Microresonators





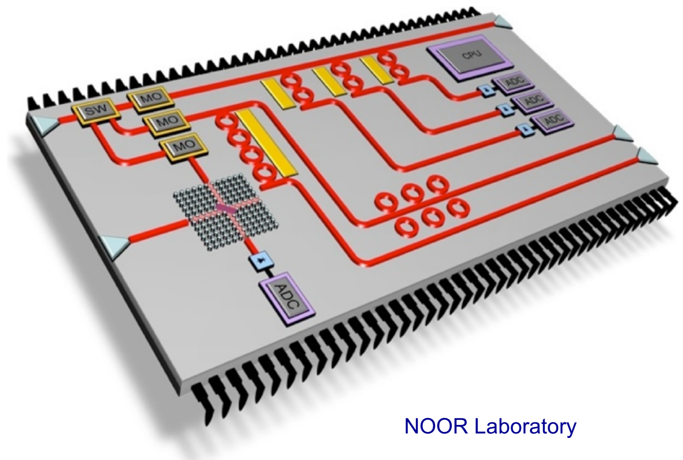




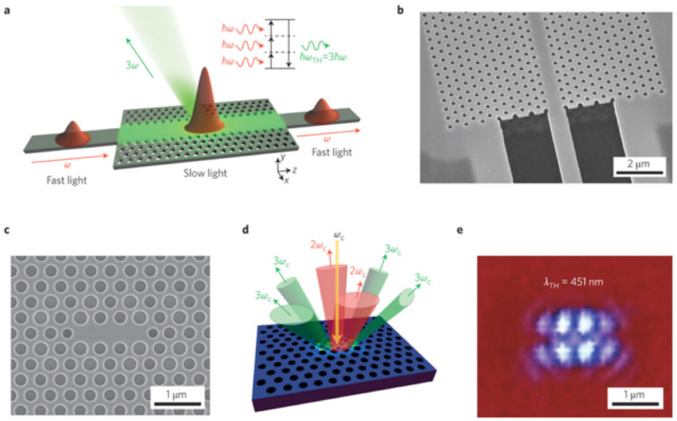
L Tong *et al*. *Nature* 426, 816 (2003)

B. Eggleton., Nat. Phot. 5, 141 (2011)

From systems to integration: The photonic circuit



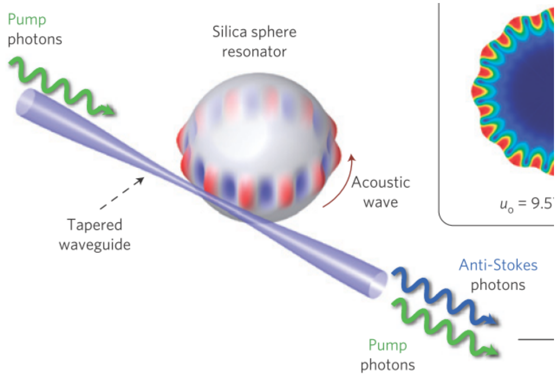
NOOR Laboratory



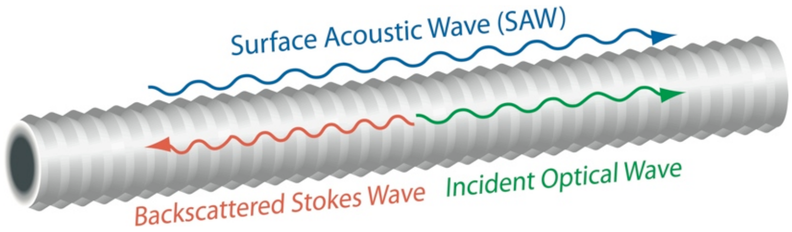
T. Krauss, Nature Nanotechnology 9, 19 (2014)

What next ?

Harnessing new nonlinear effects (opto-mechanical & radiation pressure)

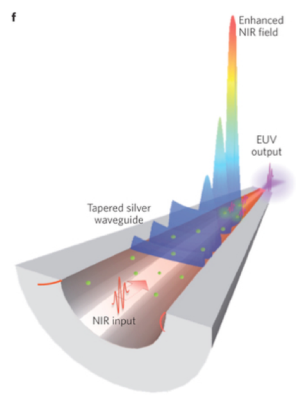


G. Bahl *et al.*, Nat. Comm., 2, 403 (2011)

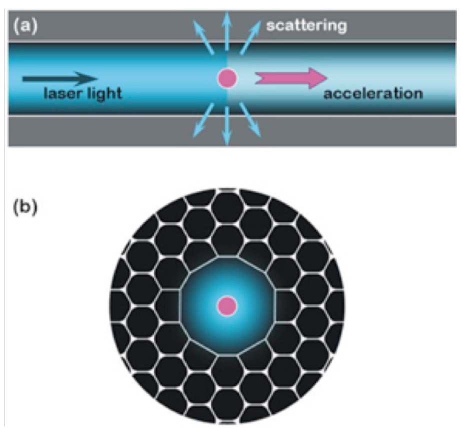


Beugnot *et al.*, Nat. Comm. (2014)

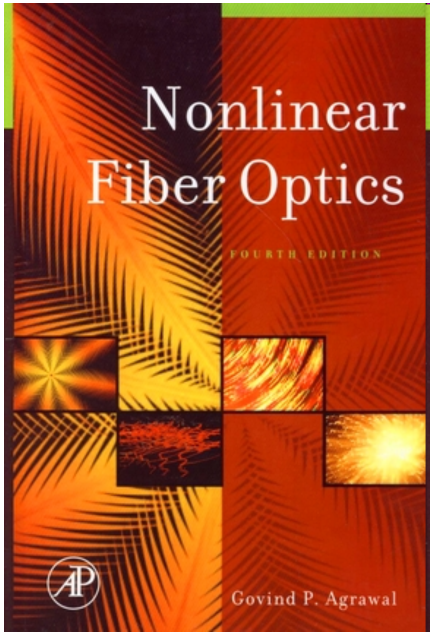
Combining nonlinear optics with nanophotonics, plasmonics and microfluidics.

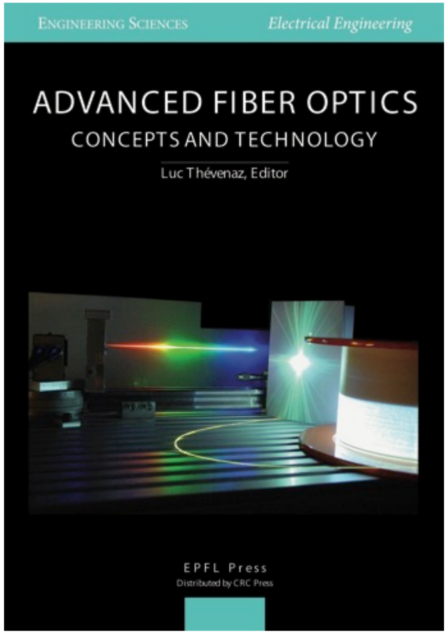


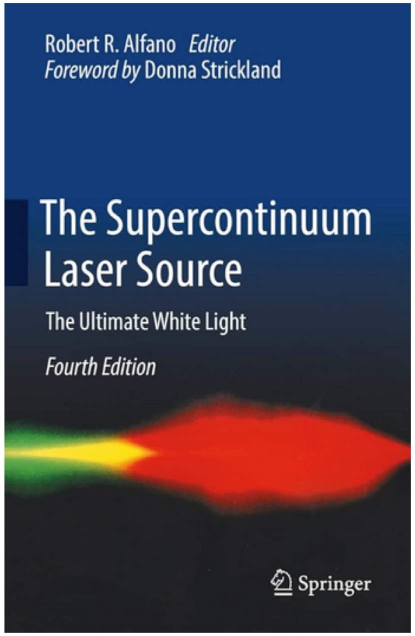
Kauranen *et al.*, Nat. Phot. 6, 737–748 (2012)

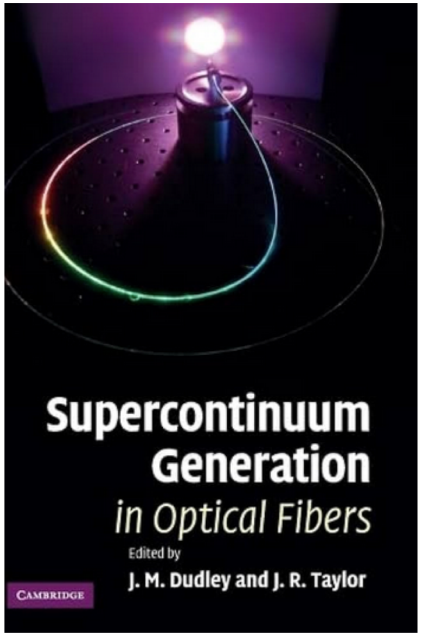


Accessible books on nonlinear fiber optics





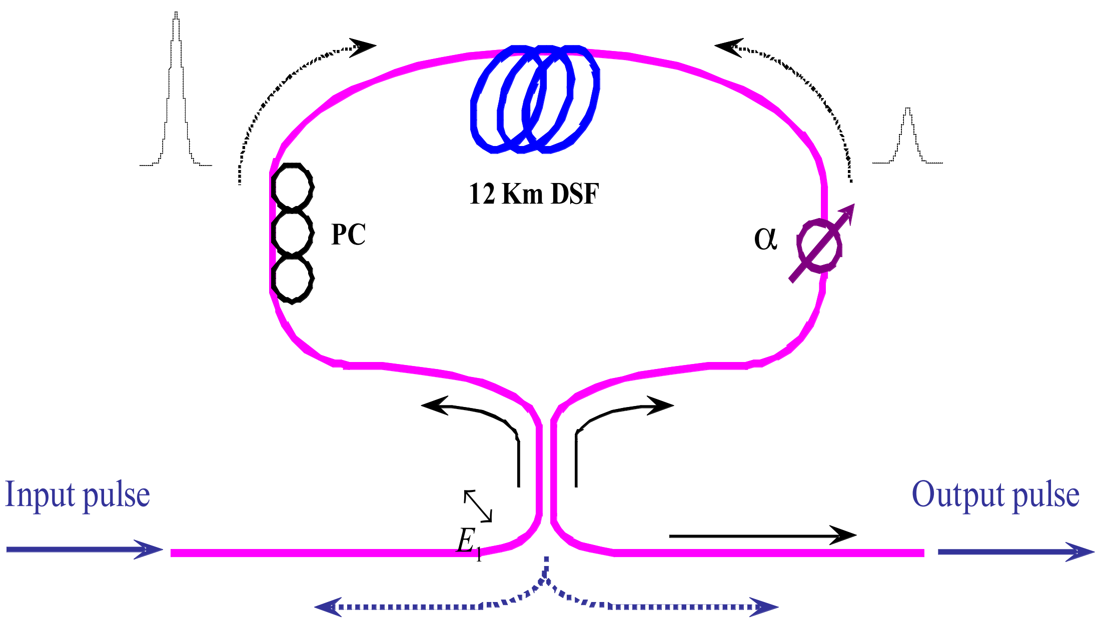




Questions ?


International Day of Light
Opening Ceremony
UNESCO, Paris, May 16, 2018

The Nonlinear Optical Loop Mirror (NOLM)

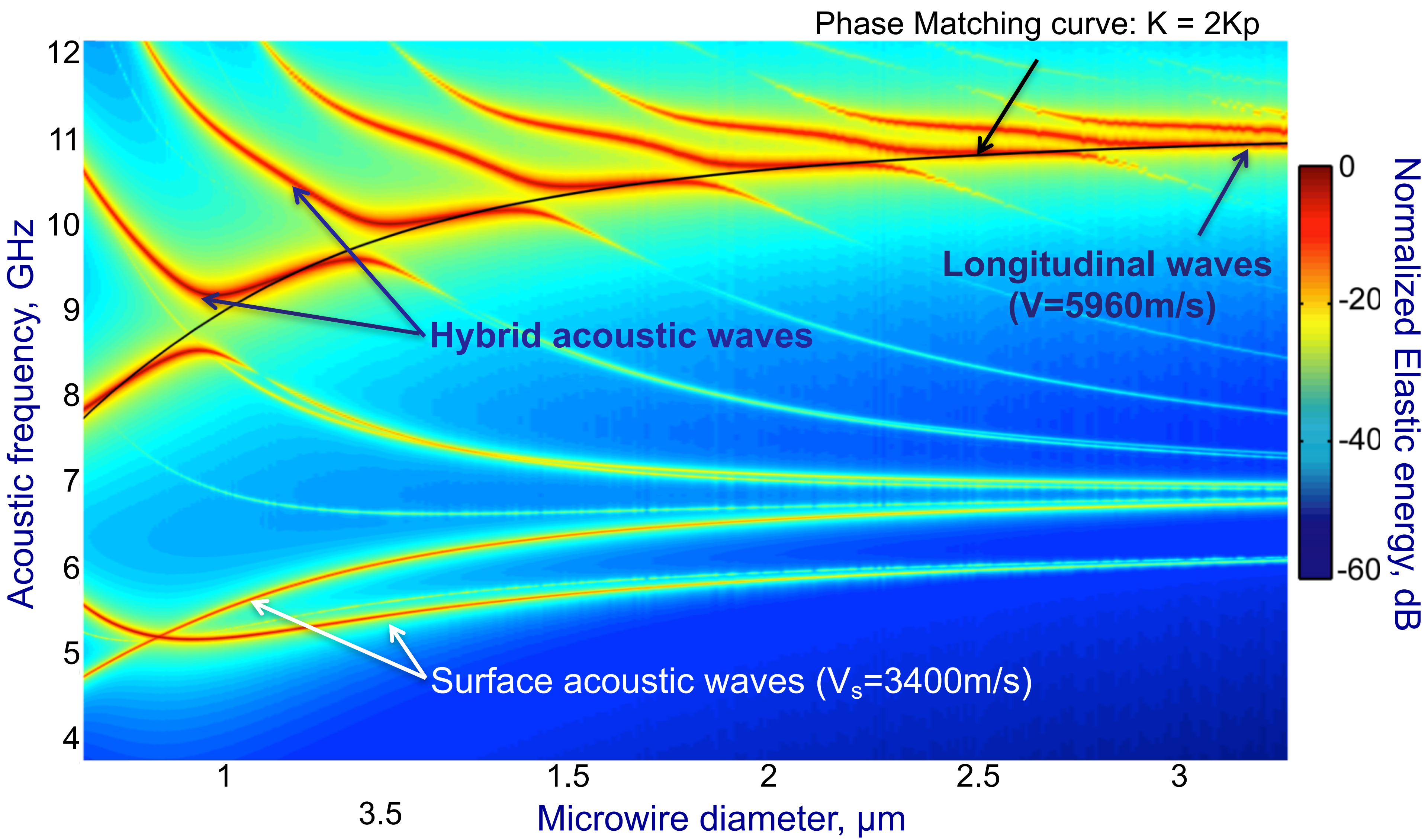


$$T = \alpha \sin^2\left(\frac{\Delta\varphi_{NL}}{2}\right)$$

$$\Delta\varphi_{NL} = \frac{\gamma P_0 L}{2}(1-\alpha)$$

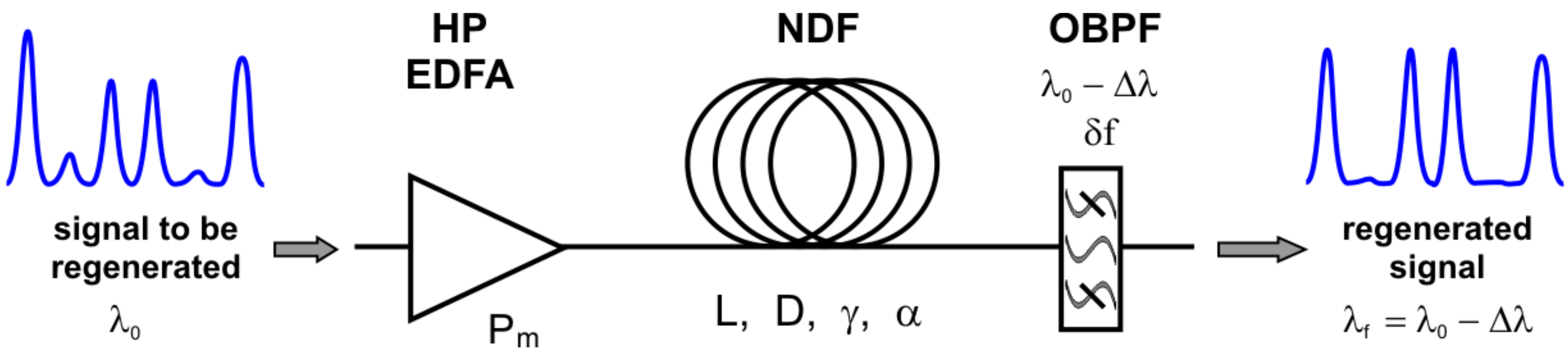
At low input powers the loop acts as a perfect mirror, and no light exits to the output port. In the high power regime, however, the refractive index of the fiber is modified by SPM, leading to pulse transmission and compression.

Brillouin scattering in a fiber taper



Applications of SPM - XPM

2R Mamyshev regenerator: pulse reshaping & saturable absorber



Mamyshev, P. V. (1998). ECOC '98, 1. pp. 475–476.